

Emerging Magnetic Order

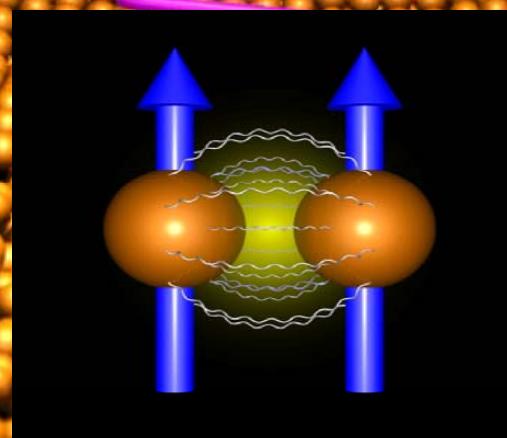


*Theo Rasing
Radboud Universiteit
Nijmegen*

Magnetism and the timescales of the exchange interaction

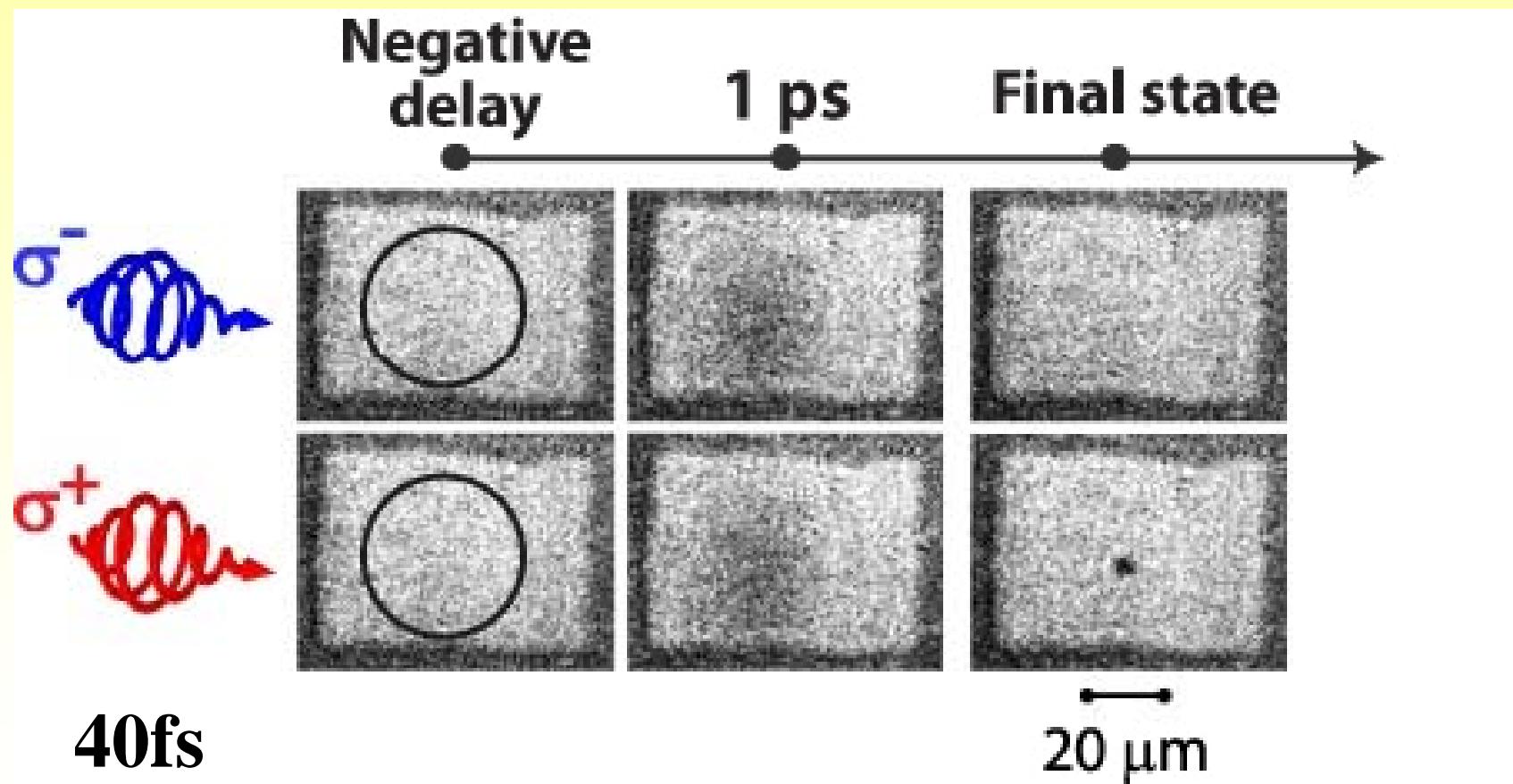
Theo Rasing

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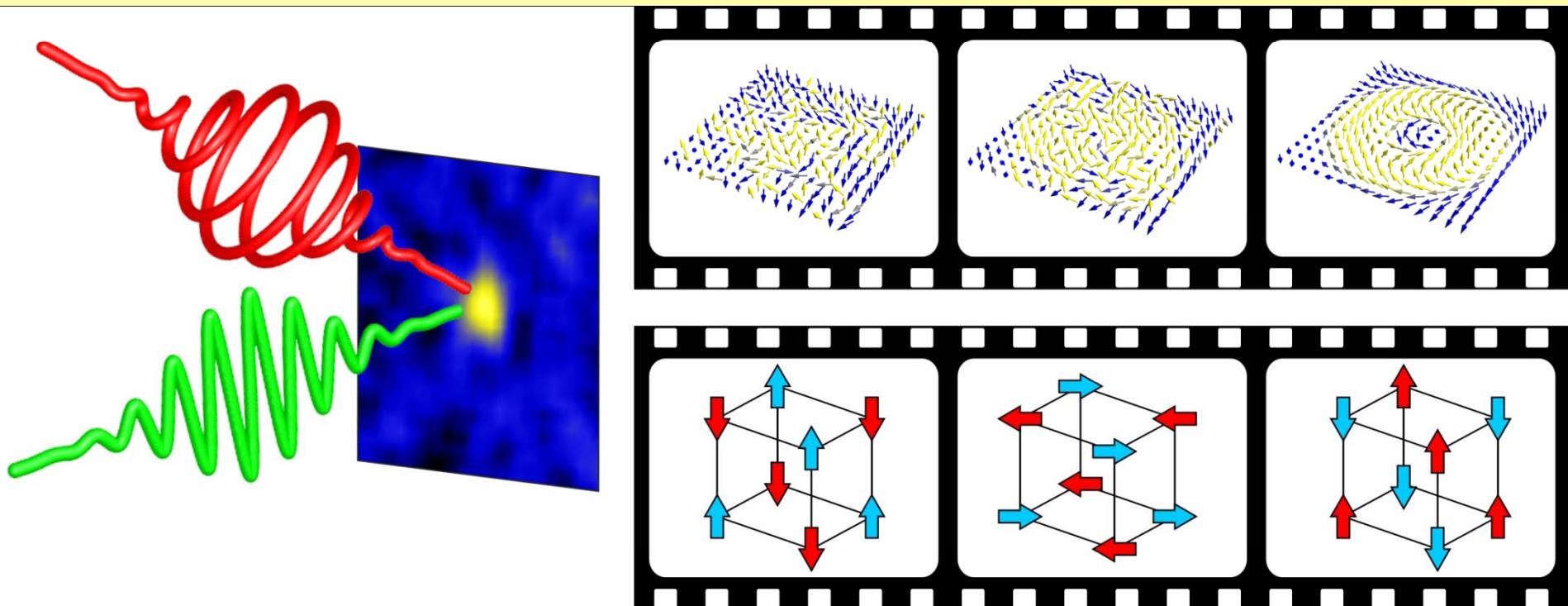


$$J_{\text{Fe-Fe}} = 1.4 \times 10^{-19} \text{ J} \sim 40 \text{ fs}$$

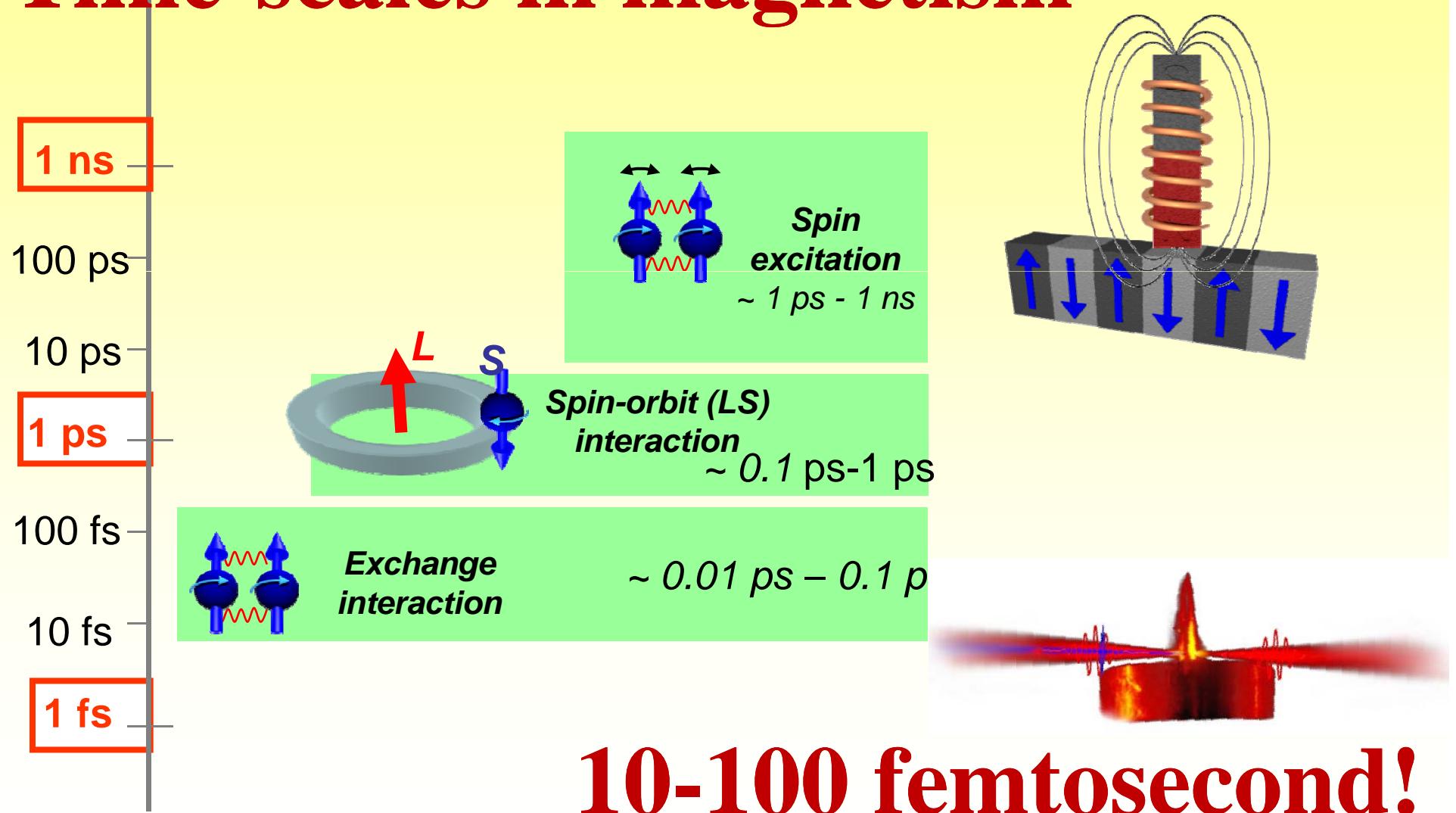
Emerging Magnetic Order



Emerging Magnetic Order



Time-scales in magnetism



Length-scales in magnetism

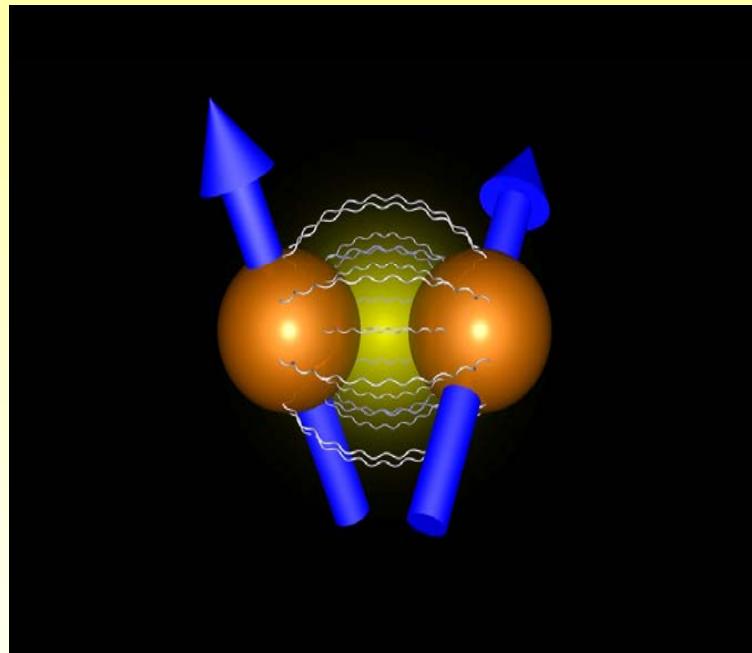
$$l = \sqrt{J/K} \sim nm$$



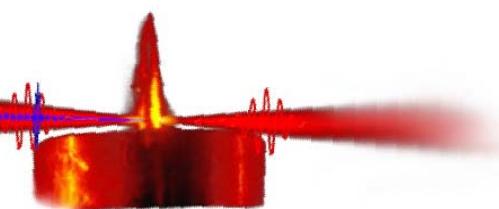
Tb+THz

nm+fs!

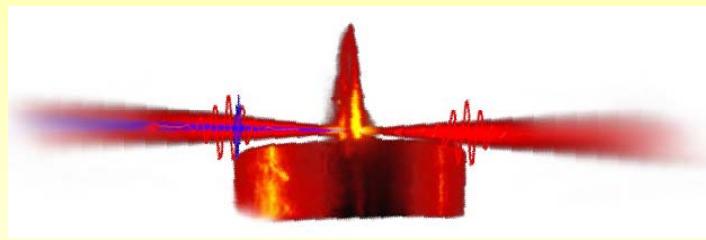
New Mechanisms?



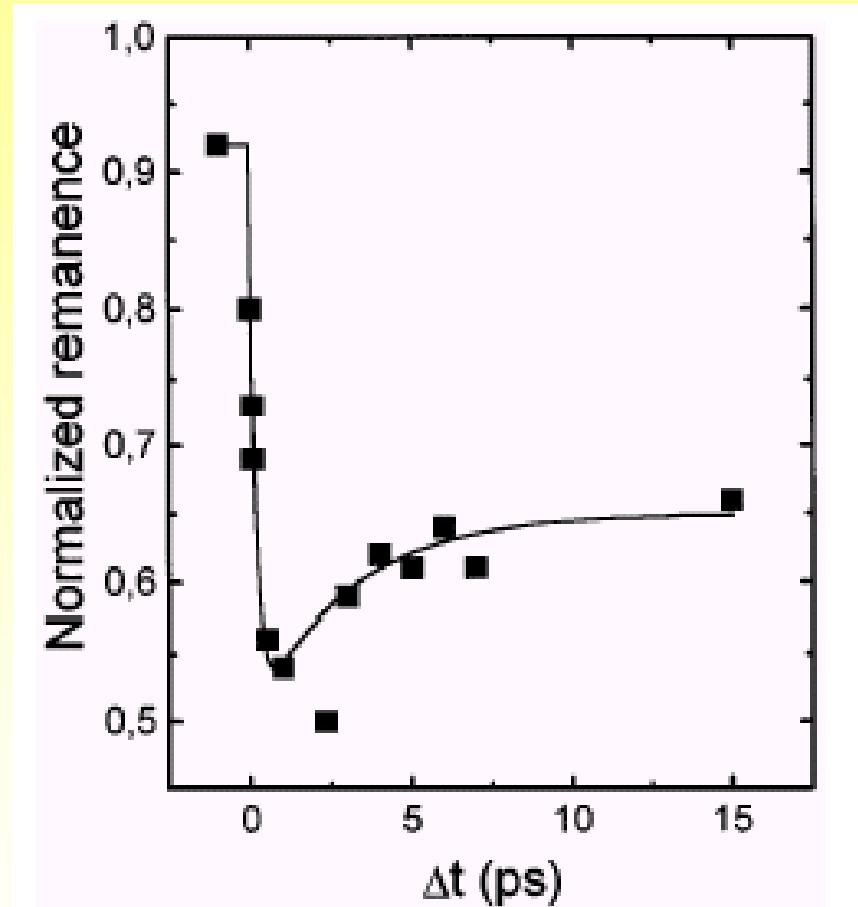
1-2eV, 10-100fs



1996 Beaurepaire et al:

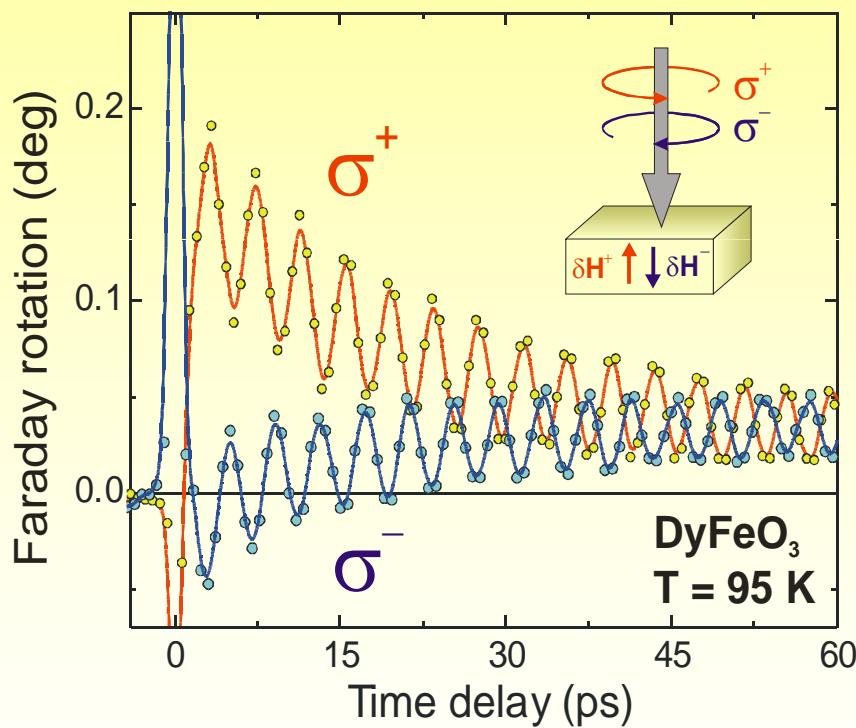


60fs



Magnetization is changed within a picosecond!!!

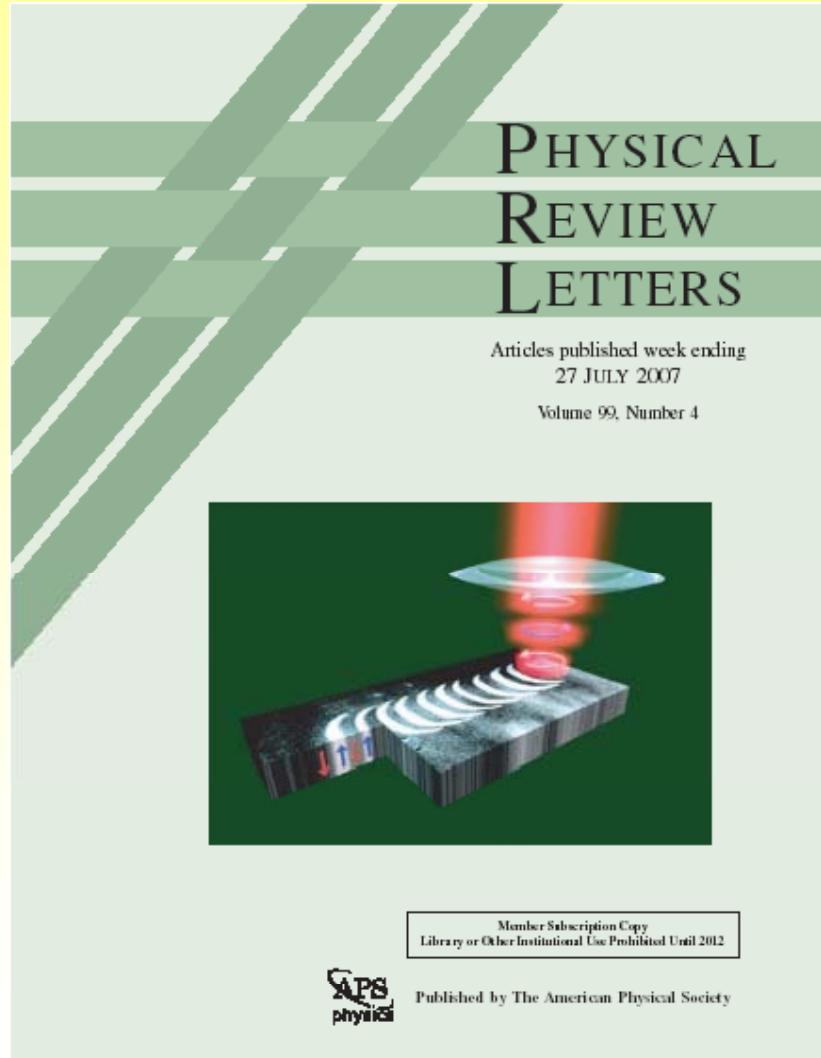
Ultrafast excitation of spins in DyFeO₃ via the Inverse Faraday effect



A. Kimel et al, *Nature* **435** 655 (2005).

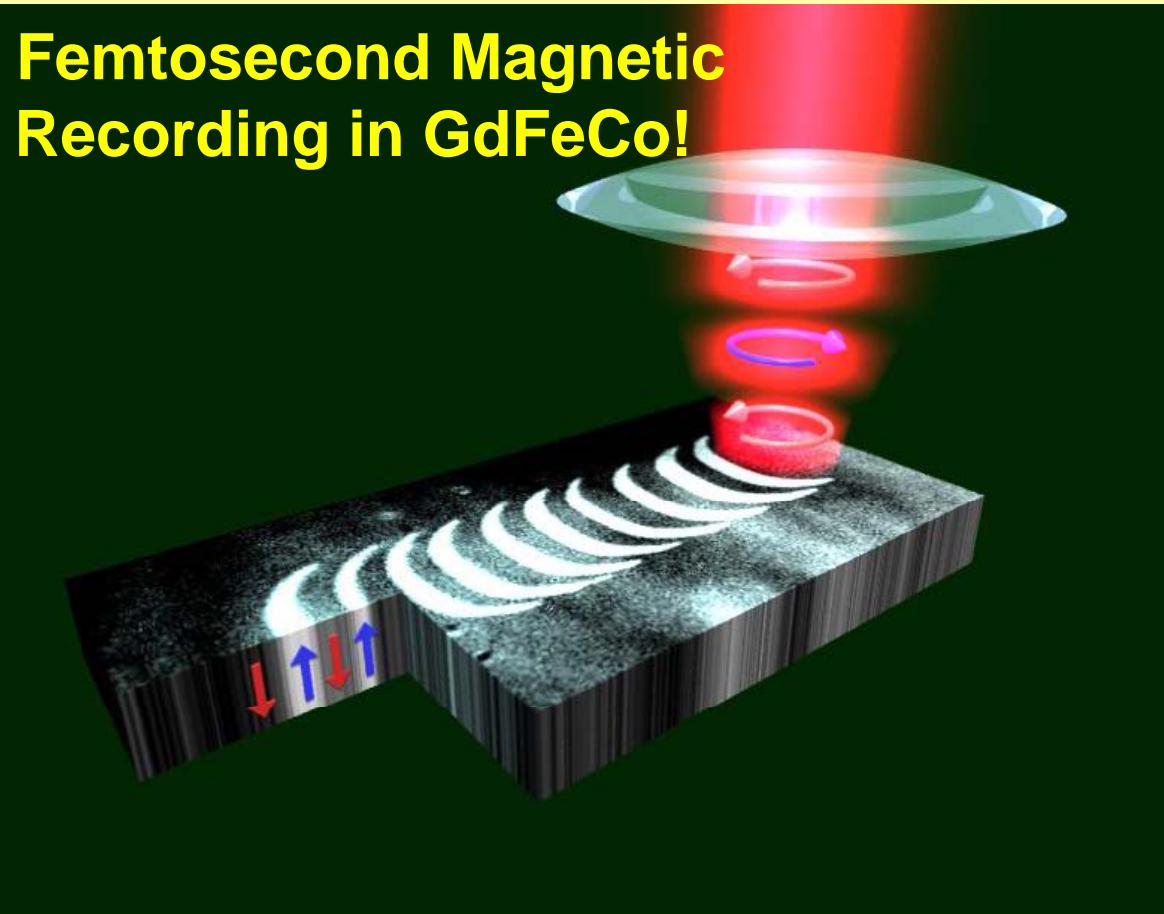
Referee A:
*“But, unfortunately,
the observed signal
is so small that it
seems
impractical to utilize
the inverse Faraday
effect for the
purpose
of ultrafast control
of magnetization in
metallic materials.”*

Ultrafast all optical recording

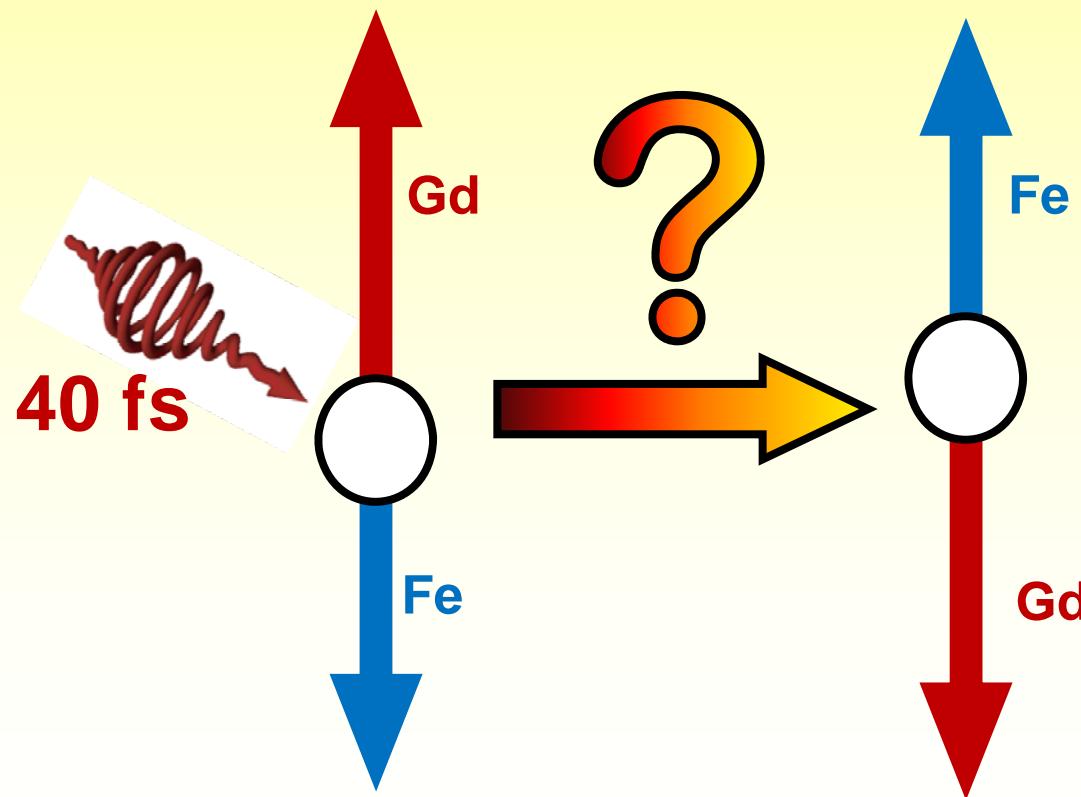


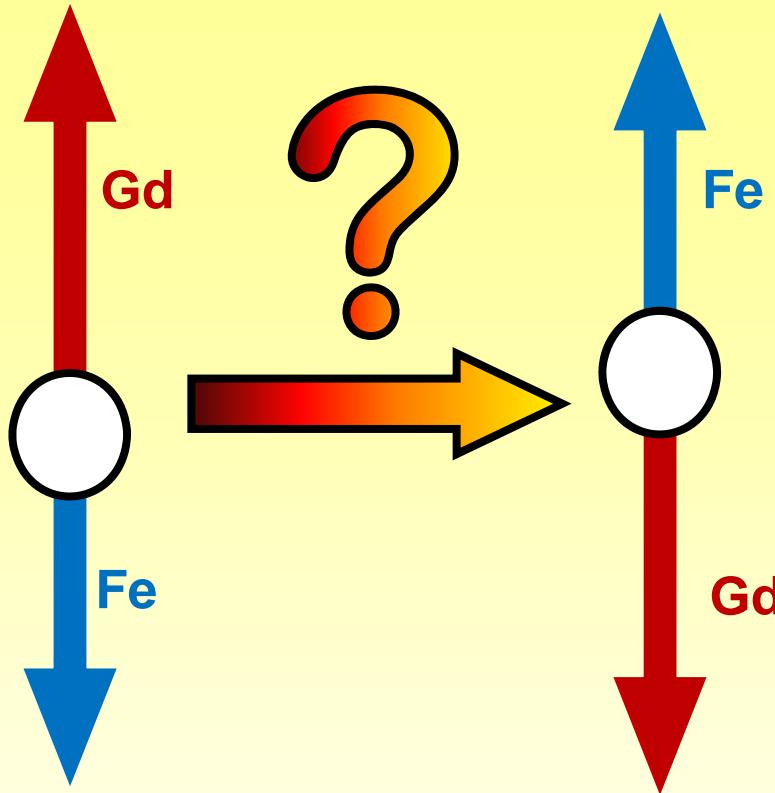
C.D. Stanciu et al., *patent #P77323PC00, PRL 99, 047601 (2007)*

GdFeCo: multi-sublattice ferrimagnet!



Ultrafast laser induced reversal in *multi*-sublattice magnets





Angular momentum transfer?
Role of temperature?
Role of sublattices?

.....

Beyond Landau-Lifshitz dynamics

ON THE THEORY OF THE DISPERSION OF MAGNETIC
PERMEABILITY IN FERROMAGNETIC BODIES.

By L. Landau and E. Lifshitz.

(Received June 3, 1935).

magnetic moment. Therefore in the presence of the field the magnetic moment would act as a free moment, i. e. would rotate around \mathbf{f} and we should have for \mathbf{s} (\cdot denotes differentiation by time) the equation

$$\dot{\mathbf{s}}/\mu_0 = [\mathbf{f}\mathbf{s}] + \lambda \left(\mathbf{f} - \frac{(\mathbf{f}\mathbf{s})\mathbf{s}}{s^2} \right). \quad (21)$$

The second term here is a vector directed from \mathbf{s} to \mathbf{f} . The constant λ is $\lambda \ll s$ in accordance with the fact that the relativistic interaction is weak. We disregard here altogether the variation of the absolute value of \mathbf{s} .

$$d\mathbf{S}/dt = -\gamma \mathbf{S} \times \mathbf{H}$$

$$\gamma = g \frac{e}{2m} = 0.28 \text{ GHz/T}$$

Typical laboratory fields $\sim 1\text{T}$
 \rightarrow Reversal time $\sim 1\text{ ns!}$

Magnetism on the timescale of exchange:

\rightarrow LONGITUDINAL spin dynamics

$$\frac{d|\mathbf{S}|}{dt} = ?$$

Onsager's relations for spin dynamics

Baryakhtar, JETP 1984

Magnetic free energy

$$F = \int \{ f(S^2) - \mathbf{S} \cdot \mathbf{H}_0 - K_u S_z^2 + \dots \} dV$$

$$\frac{dF}{dt} = - \int \frac{d\mathbf{S}(\mathbf{r}, t)}{dt} \cdot \mathbf{H}(\mathbf{r}, t) dV$$

Generalized force
 $\mathbf{H}(\mathbf{r}, t) = -\delta F / \delta \mathbf{S}(\mathbf{r}, t)$

Generalized position
 $\mathbf{S}(\mathbf{r}, t)$ angular momentum

Symmetry of kinetic coefficients

$$\begin{aligned} dS_i(\mathbf{r}', t)/dt &= \int \lambda_{ik}(\mathbf{r}' - \mathbf{r}, \mathbf{S}(\mathbf{r})) H_k(\mathbf{r}', t) dV \\ &= \lambda_{ik}(\mathbf{S}) H_k(\mathbf{r}, t) + \lambda_{ik,lm} \partial^2 H_k / \partial r_l \partial r_m + \dots \end{aligned}$$

$$\lambda_{ik}(\mathbf{S}) = \lambda_{ki}(-\mathbf{S})$$

$$\begin{aligned} \lambda_{ik}^s(\mathbf{S}) &= (\lambda_{ik} + \lambda_{ki})/2 \\ \lambda_{ik}^a(\mathbf{S}) &= (\lambda_{ik} - \lambda_{ki})/2 \end{aligned}$$

$$d\mathbf{S}/dt = \mathbf{S} \times \mathbf{H} + \lambda^s \mathbf{H} + \dots$$

Anti-symmetric Symmetric
precession longitudinal

More general: **multiple** sublattices

Rate of change free energy

$$\frac{dF}{dt} = - \int \left\{ \frac{d\mathbf{S}_1}{dt} \cdot \mathbf{H}_1 + \frac{d\mathbf{S}_2}{dt} \cdot \mathbf{H}_2 \right\} dV$$

Longitudinal dynamics to lowest order

$$\frac{d}{dt} \begin{pmatrix} S_1 \\ S_2 \end{pmatrix} = \begin{pmatrix} \lambda_{11}^s & \lambda_{12}^s \\ \lambda_{21}^s & \lambda_{22}^s \end{pmatrix} \begin{pmatrix} H_1 \\ H_2 \end{pmatrix}$$

$$\begin{pmatrix} \lambda_1 & 0 \\ 0 & \lambda_2 \end{pmatrix} + \begin{pmatrix} \lambda_e & -\lambda_e \\ -\lambda_e & \lambda_e \end{pmatrix}$$

relativistic exchange

$$dS_1/dt + dS_2/dt = 0$$

J.H. Mentink *et al.*, *Phys. Rev. Lett.* (2012)

More general: multiple sublattices

$$\begin{aligned} dS_1/dt &= \lambda_e(H_1 - H_2) + \lambda_1 H_1 \\ dS_2/dt &= -\lambda_e(H_1 - H_2) + \lambda_2 H_2 \end{aligned}$$

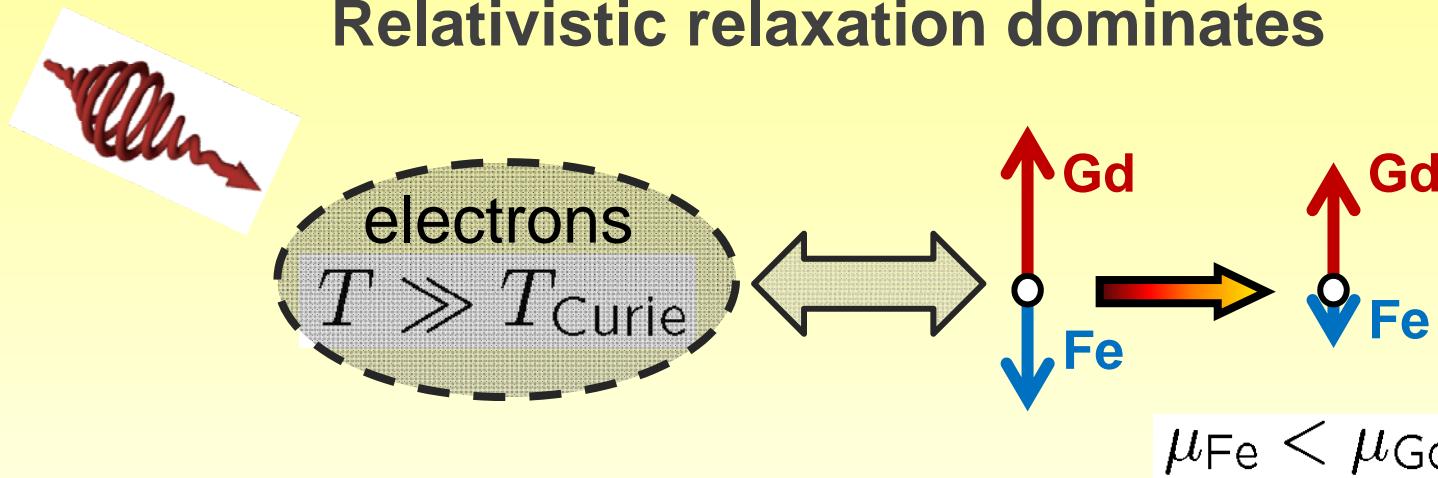
$$F = \int dV \left\{ f_1(S_1^2) + f_2(S_2^2) - J_{12}(S_1 \cdot S_2) \right\}$$

formation sublattice mutual alignment
spin moment sublattices

Temperature dominated: $T \gg T_C$, $t \sim 100\text{fs}$

Temperature dominated regime

Relativistic relaxation dominates



Distinct dynamics Gd and Fe

$$\tau_{\text{Fe}} < \tau_{\text{Gd}}$$

Bloch relaxation

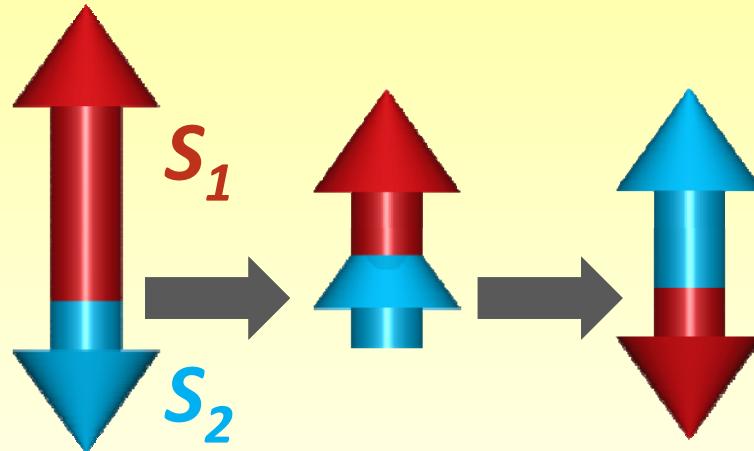
$$\frac{dM_i}{dt} = -\frac{M_i - M_i^{(0)}}{\tau_i}$$

$$\tau_i = \frac{\mu_i}{2\alpha\gamma k_B T}$$

Brown Phys. Rev. 1963
Kubo et al. Prog. Theor. Phys. Suppl. 1970

Exchange dominated: $T < T_C$, $t \sim 1\text{ps}$

Conservation total angular momentum



$$\begin{aligned} dS_1/dt &= \lambda_e(H_1 - H_2) + \cancel{\lambda_1 H_1} \\ dS_2/dt &= -\lambda_e(H_1 - H_2) + \cancel{\lambda_2 H_2} \end{aligned}$$

$$dS_1/dt = -dS_2/dt$$

What if $S_2=0$?

$$dS_2/dt|_{S_2=0} > 0 \Leftrightarrow$$

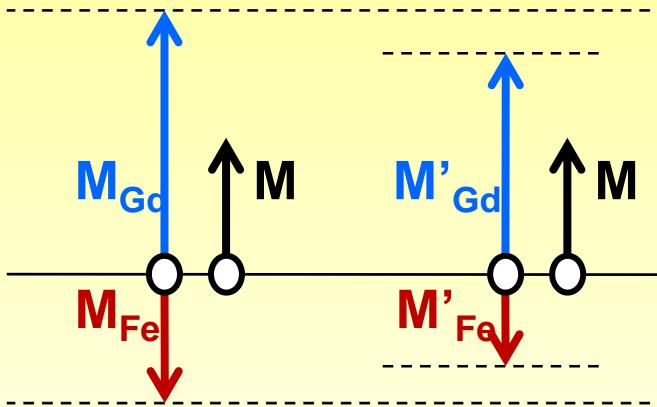
$$(S_1^2 - \bar{S}_1^2)/\chi_1 > |J_{12}|(1 + \lambda_2/\lambda_e) > 0$$

$$f_i(S_i^2) = (S_i^2 - \bar{S}_i^2)^2/4\chi_i$$

Nonequilibrium $S_1^2 \gg \bar{S}_1^2$
Equilibrium value

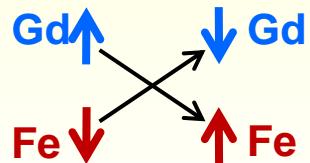
Ground state AFM, transient FM!

Exchange dominated regime



spin-spin interactions

conserve total magnetization \mathbf{M}



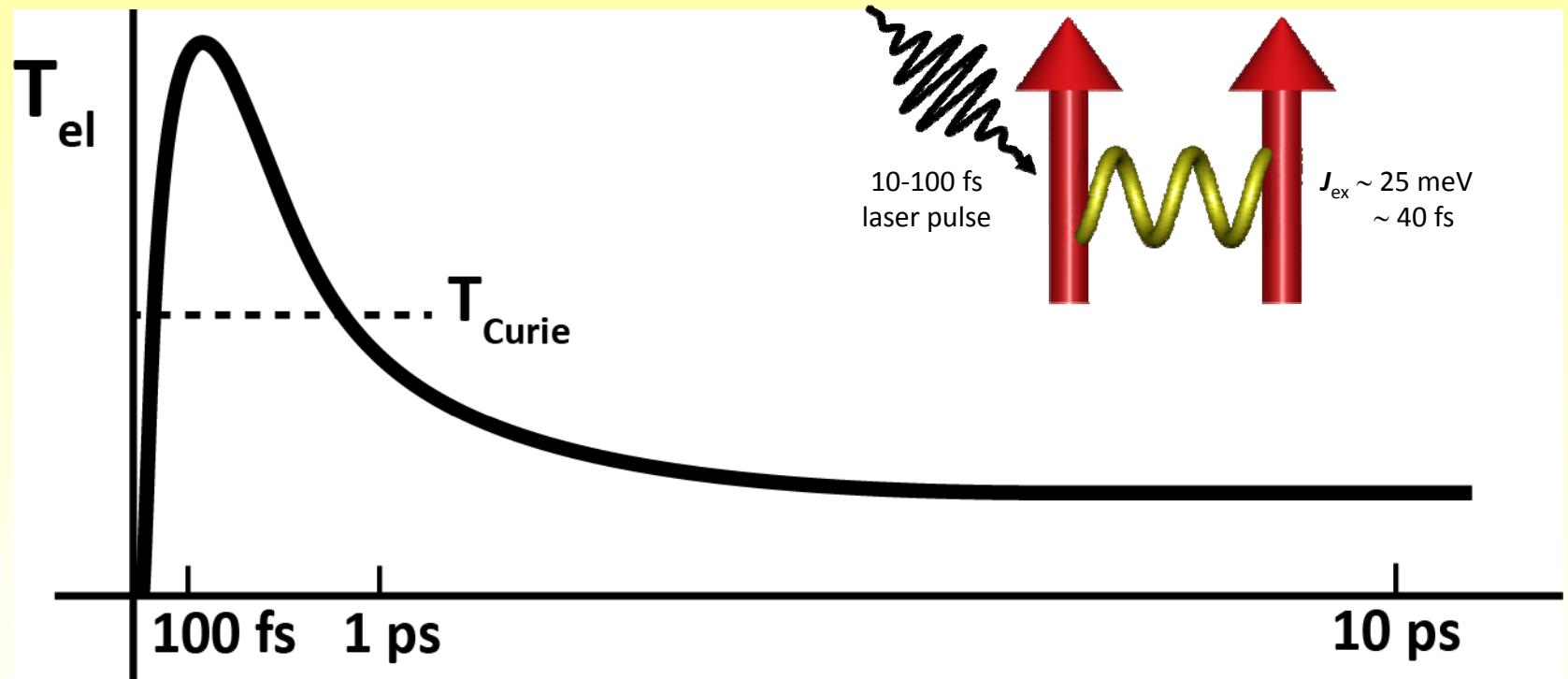
$$\frac{dM_{Fe}}{dt} = -\frac{dM_{Gd}}{dt}$$

Transfer of magnetization between sublattices!

Baryakhtar Fiz. Nizk. Temp. 1985

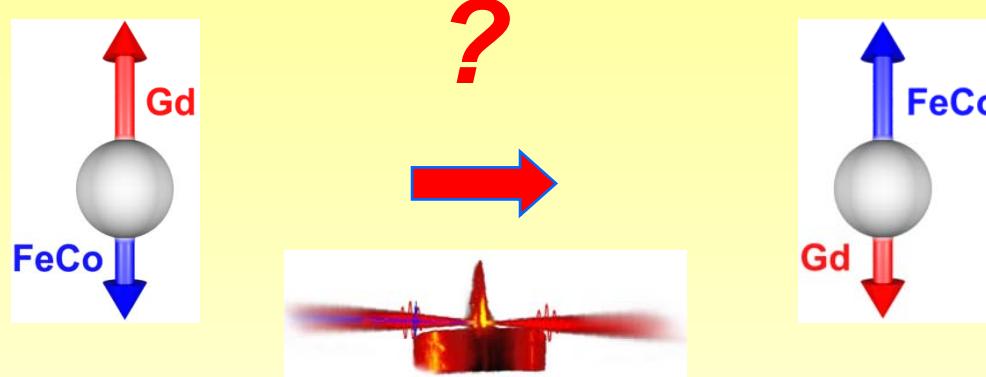
Laser-induced spin dynamics

Ultrafast heating of electrons



Access to both temperature and exchange dominated regime!

Element specific view

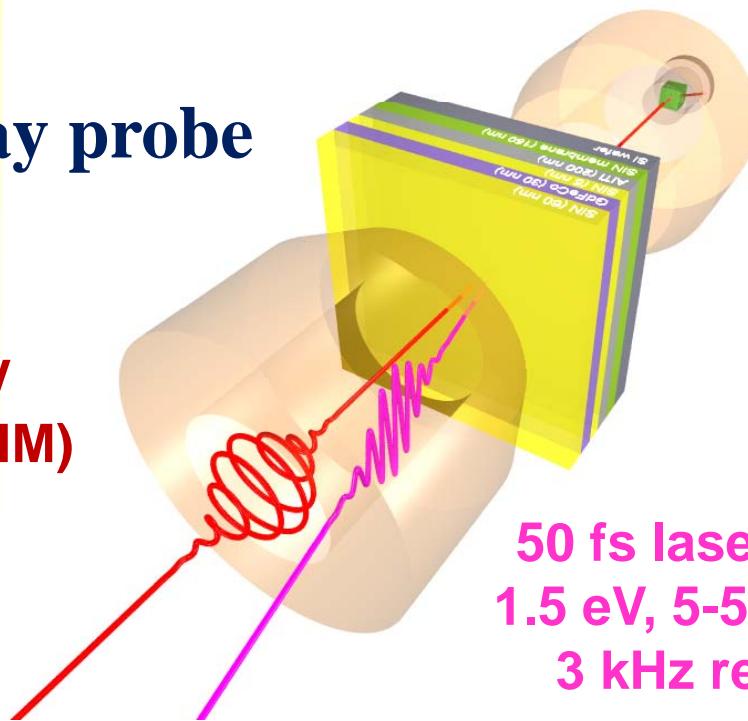


Laser pump – X-ray probe

BESSY

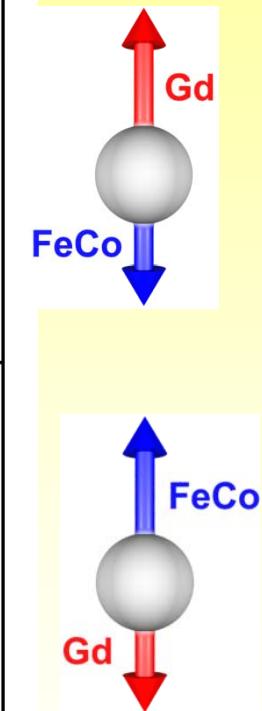
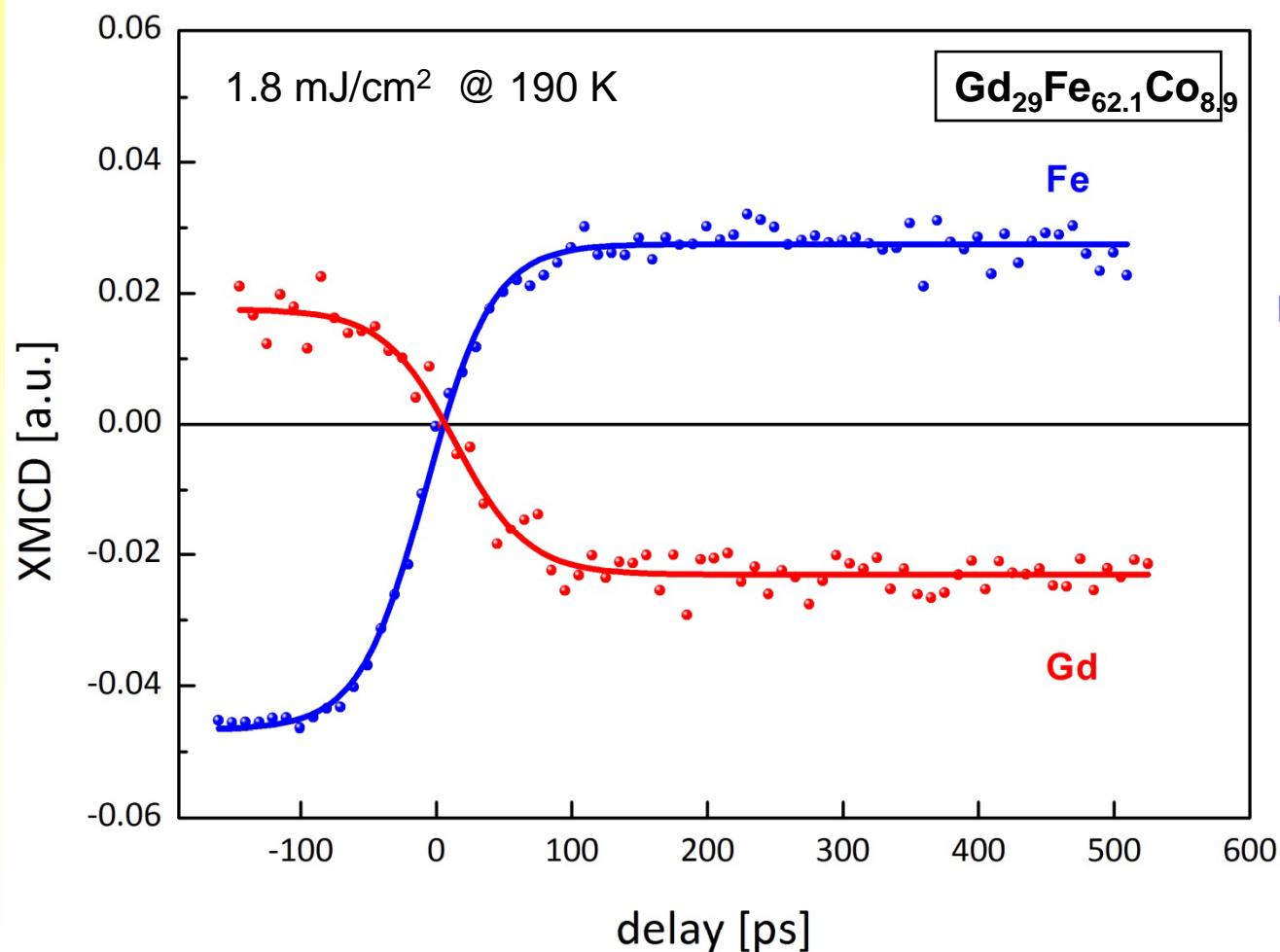
X-rays
400-1400 eV
10-50 ps (FWHM)

50 fs laser pulses
1.5 eV, 5-50 mJ/cm²
3 kHz rep. rate



TR-XMCD @BESSY

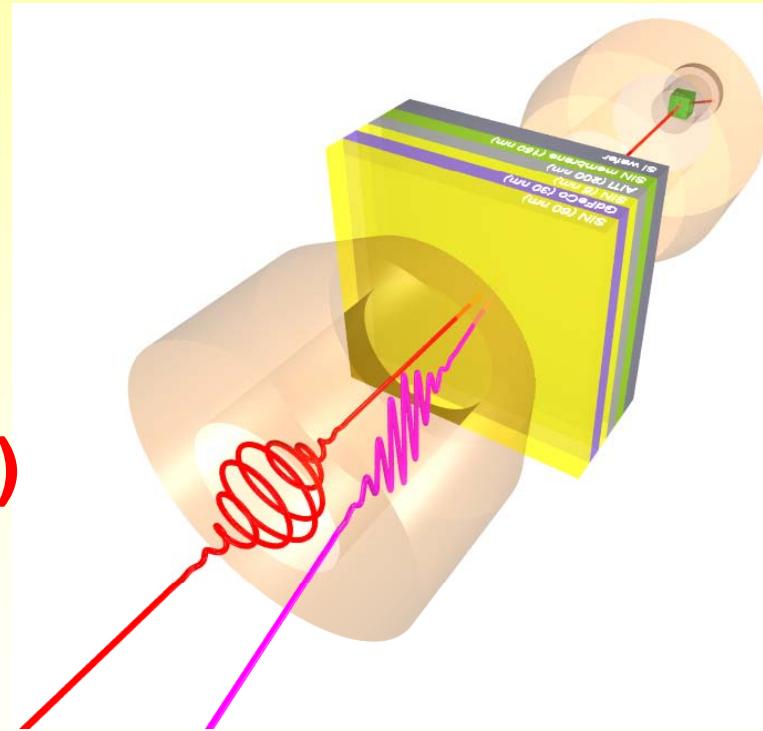
» Fe and Gd sub-lattices switch simultaneously within 50 ps



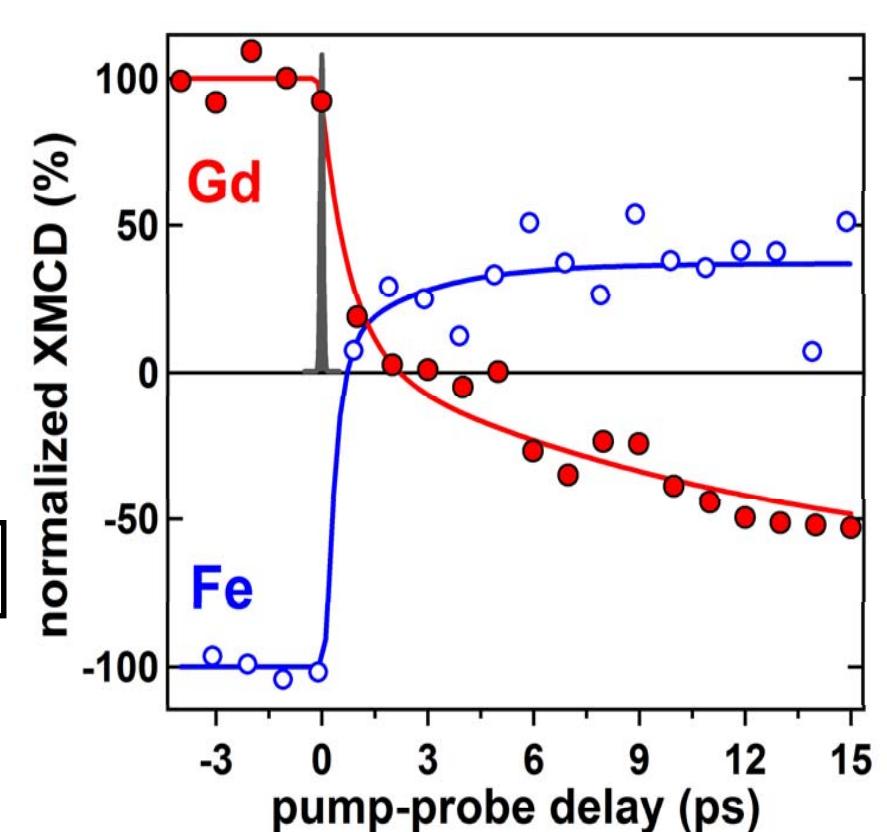
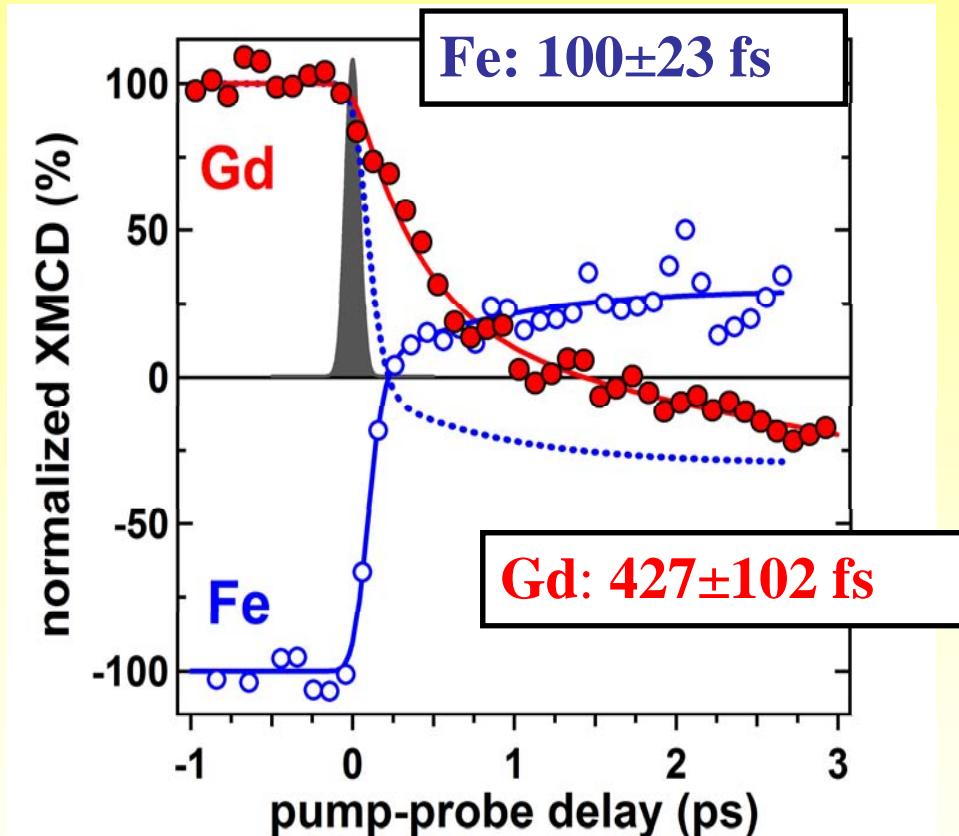
Can't we look faster?

FEMTO-SLICING!

X-rays
400-1400 eV
100 fs (FWHM)



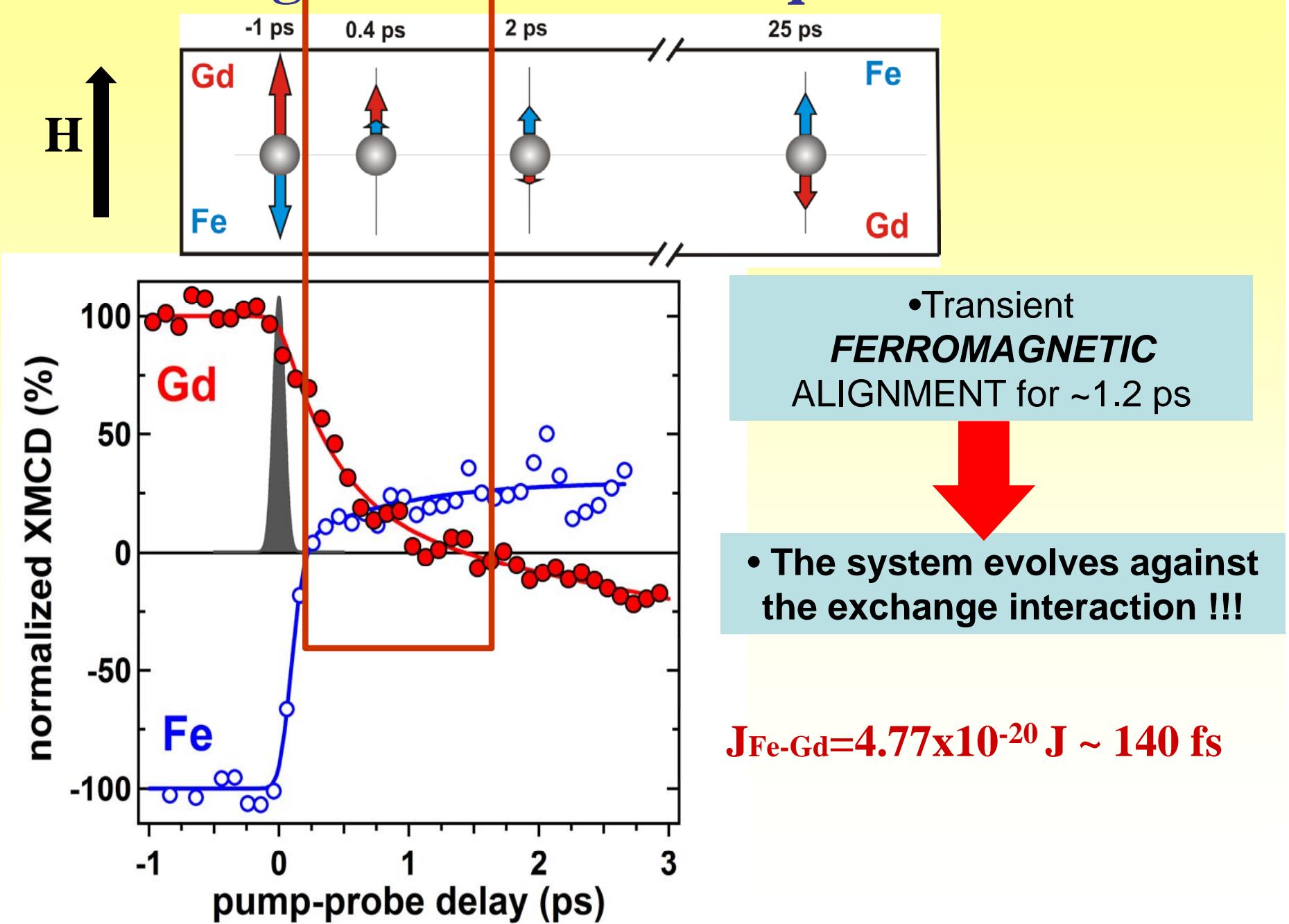
Element-specific magnetization reversal: 100 fs time resolution

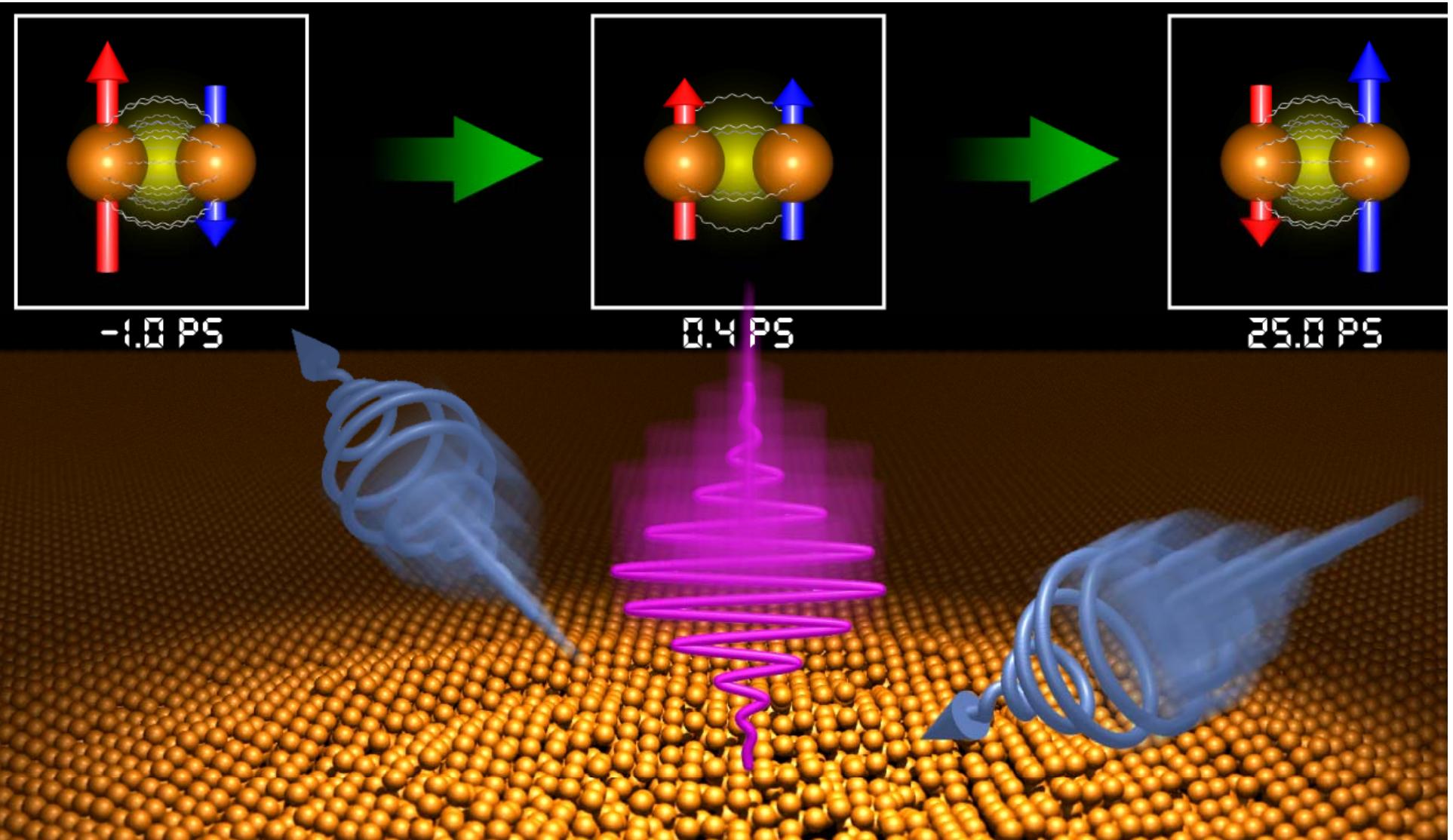


- Different magnetization switching dynamics at Fe and Gd sites !!!

Mentink: $\tau_i = \frac{\mu_i}{2\alpha\gamma k_B T}$

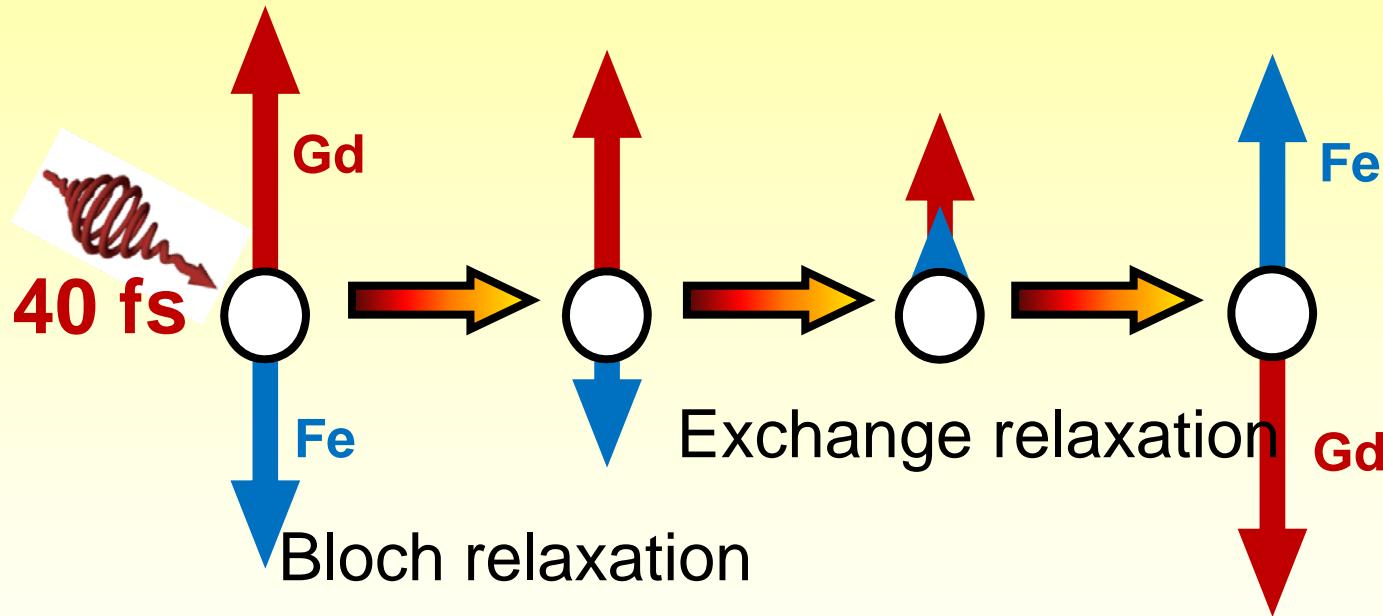
Switching via a novel non-equilibrium state





Ultrafast reversal via transient ferromagnetic phase, I. Radu et al, Nature 472, 205 (2011)

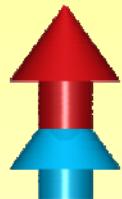
Ultrafast magnetism in multi-sublattice magnets



What about multi-sublattice ferromagnets?

Prediction: ultrafast demagnetization

FM coupling

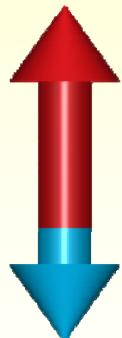


$$dS_1/dt = -\lambda_1 S_1/\chi_1 - \lambda_e(S_1/\chi_1 - S_2/\chi_2)$$

positive/negative

- Coupling makes one sublattice **faster**, other **slower**

AFM coupling



$$dS_1/dt = -\lambda_1 S_1/\chi_1 - \lambda_e(S_1/\chi_1 + |S_2|/\chi_2)$$

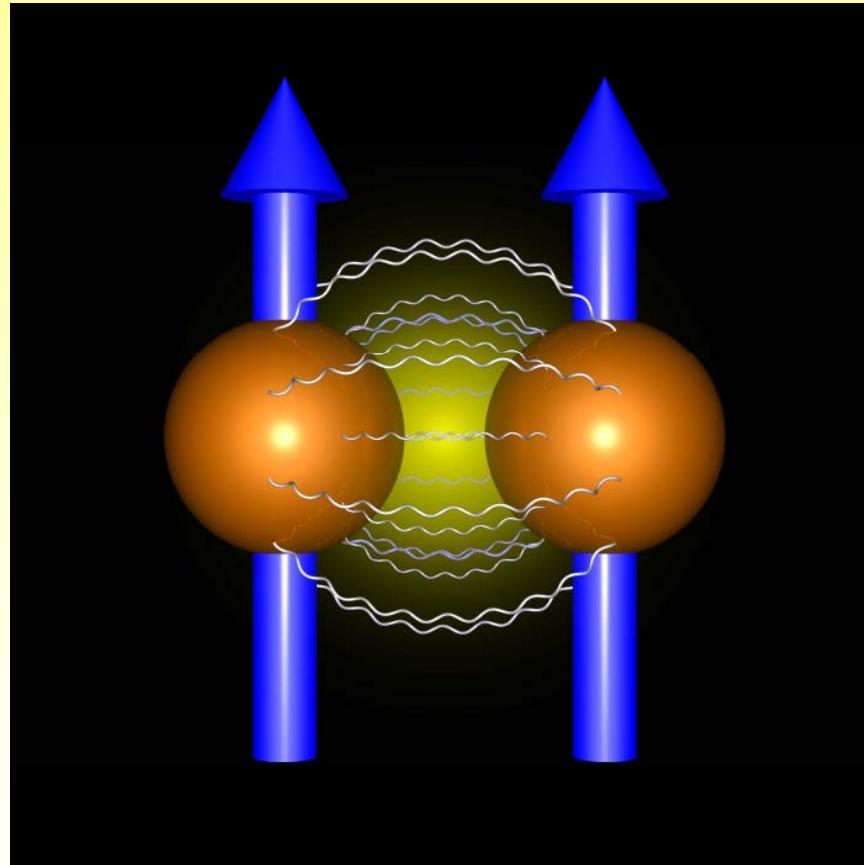
always positive

- Both sublattices **faster** than in uncoupled case!

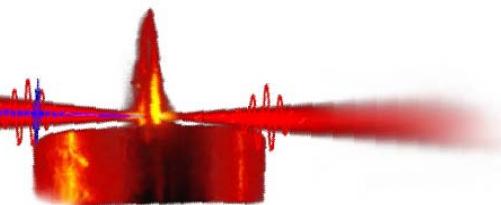
J.H. Mentink *et al.*, *Phys. Rev. Lett.* (2012)

FeNi

$$J_{\text{Fe-Ni}} = 1.4 \times 10^{-19} \text{ J} \sim 40 \text{ fs}$$



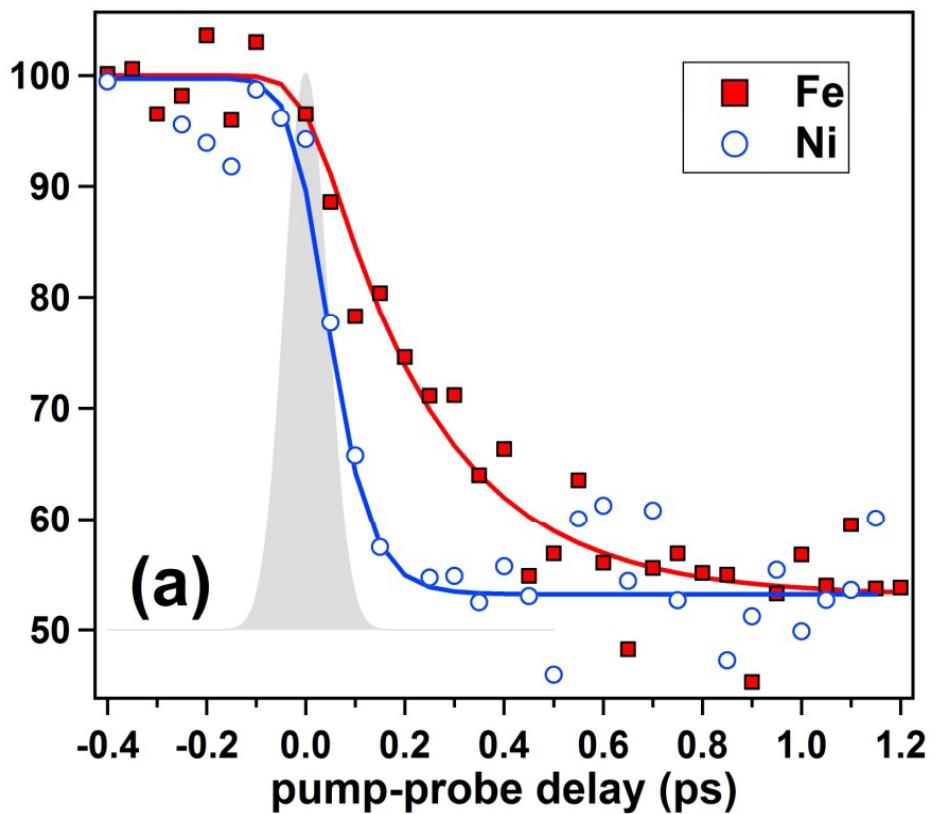
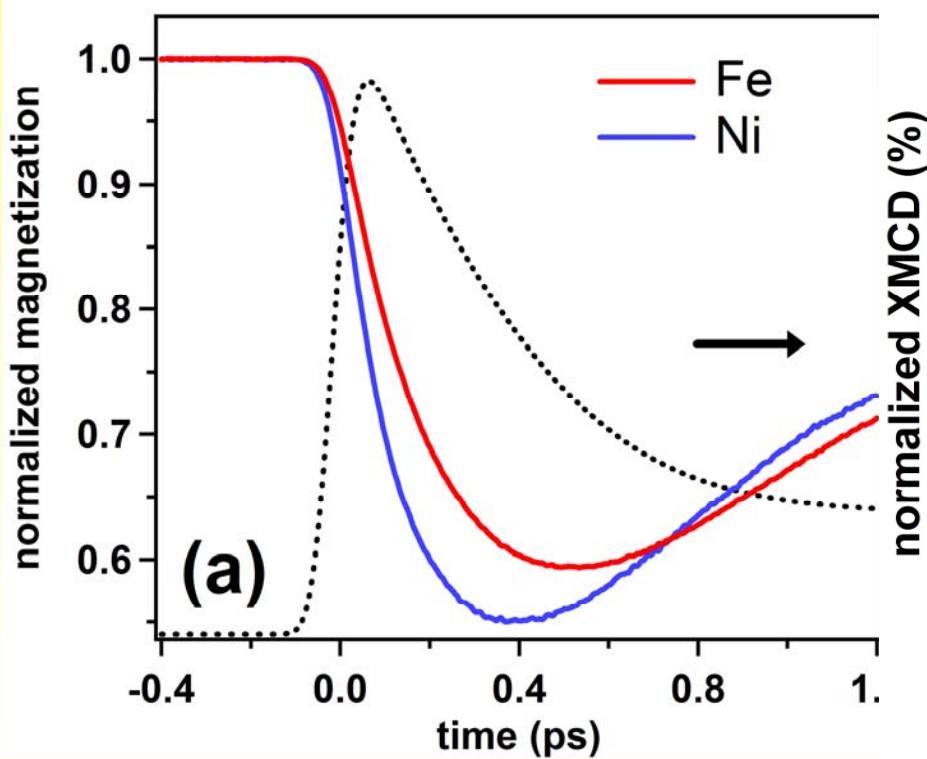
1-2eV, 10-100fs



P

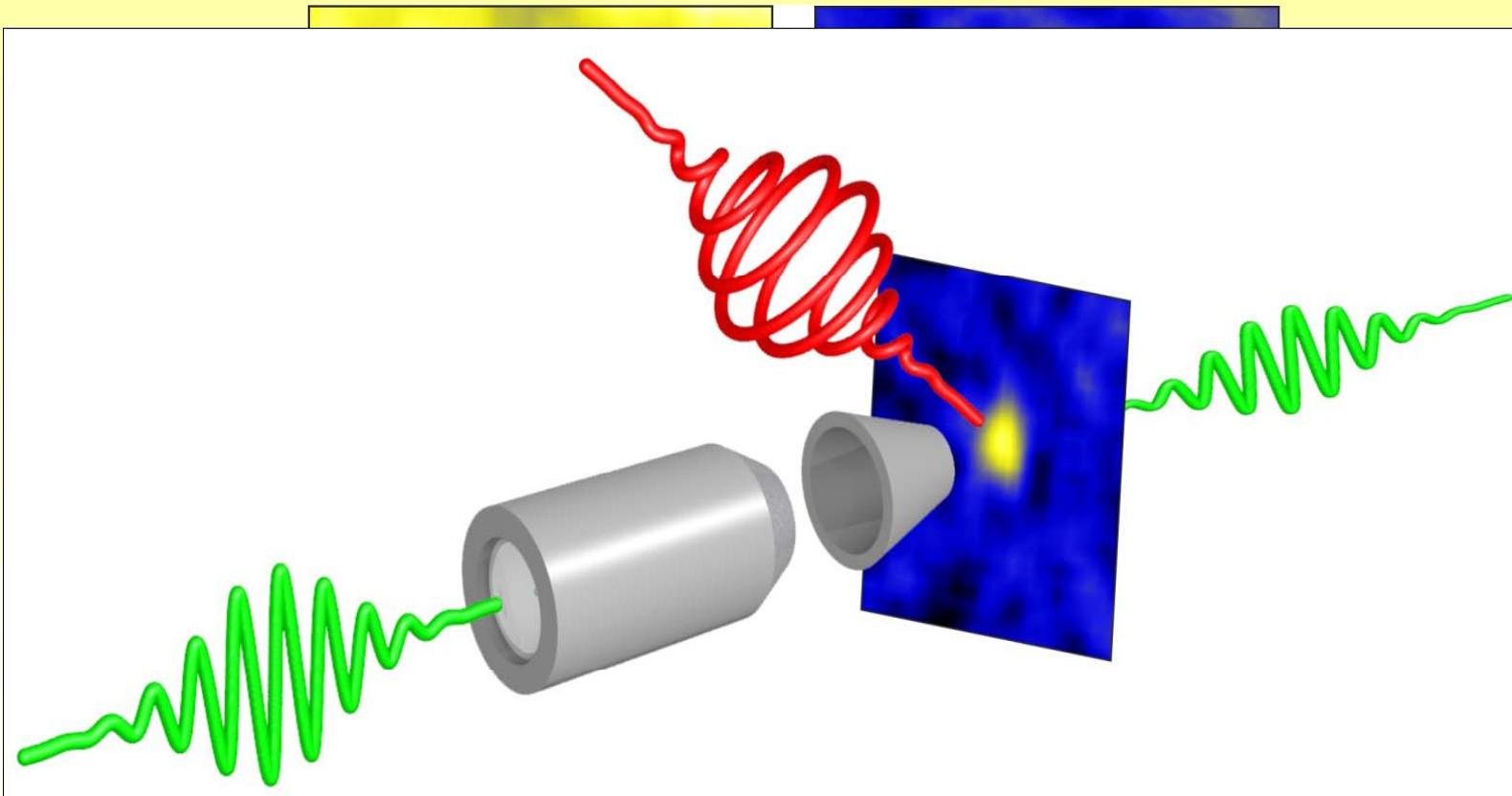
simulations

experiment



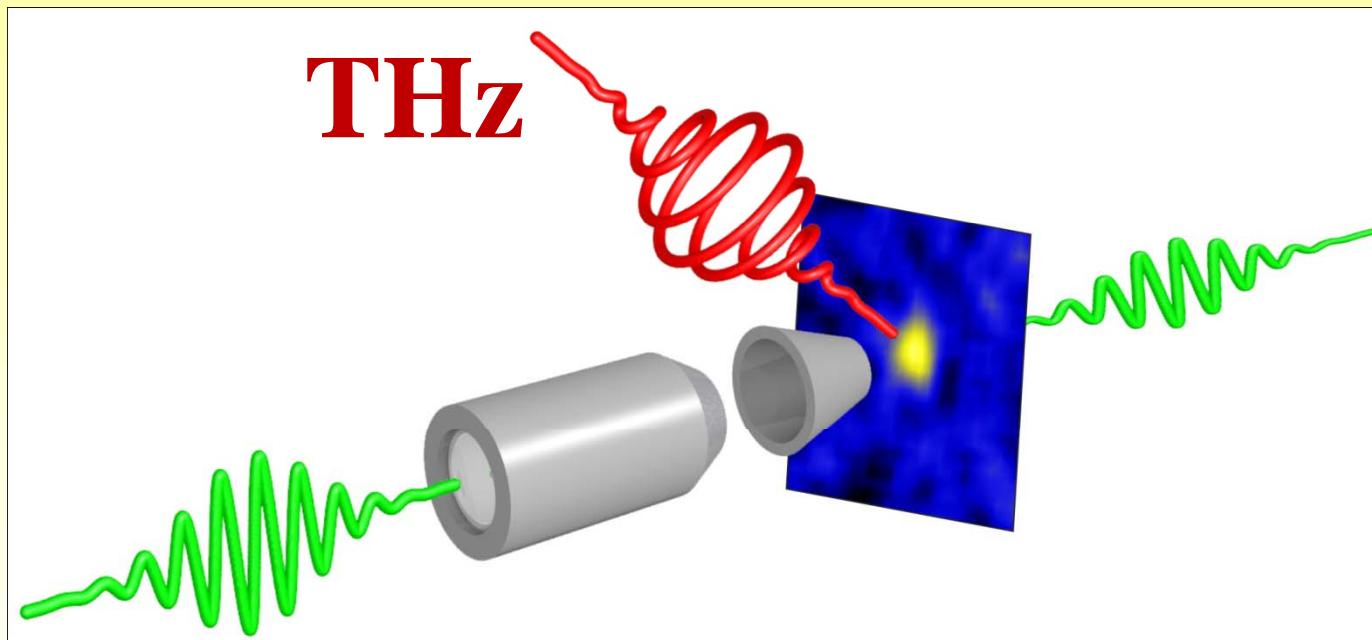
Fe slower than Ni

Perspectives



Skyrmions?

Perspectives



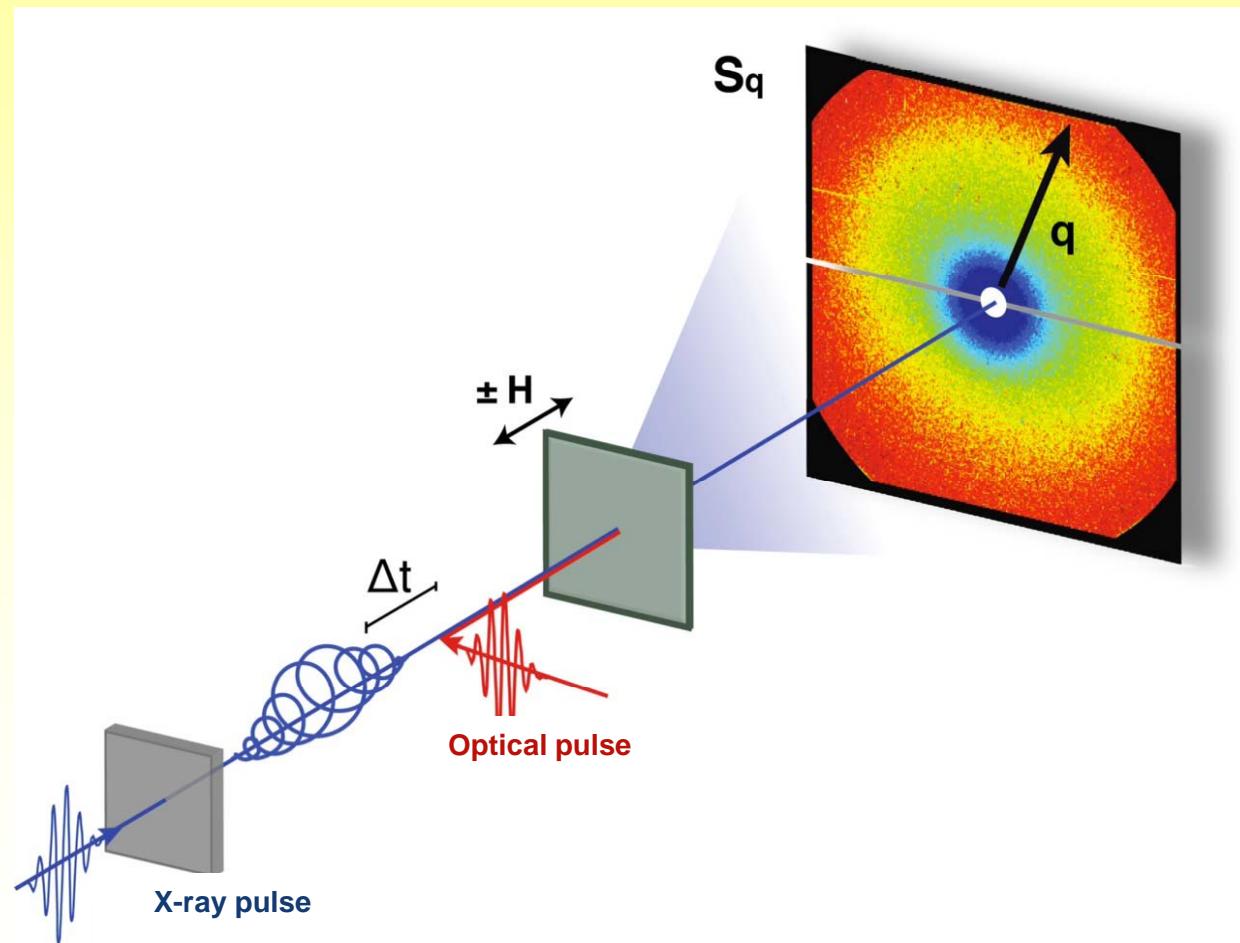
Perspectives

$J = J(t)$ \rightarrow $H(t)$ \rightarrow



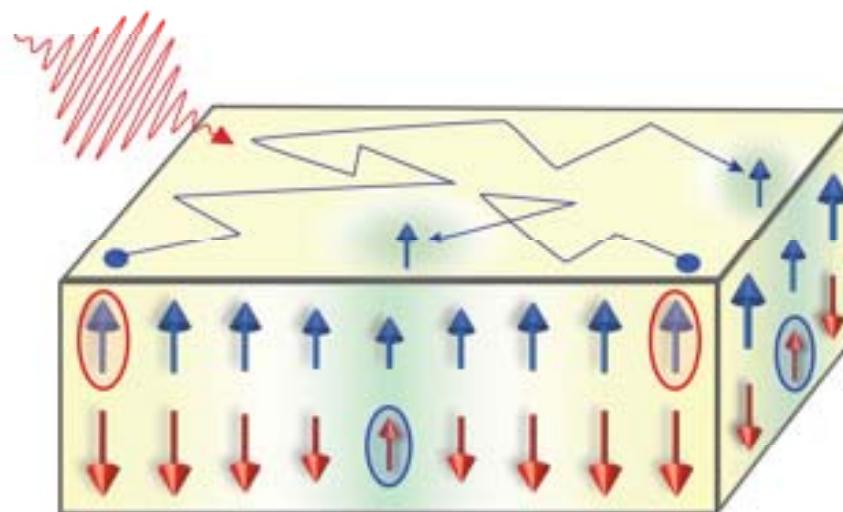
Xrays+XMCD+nm+fs!

First LCLS data



Supersonic transport of angular momentum!

(~nm/fs)



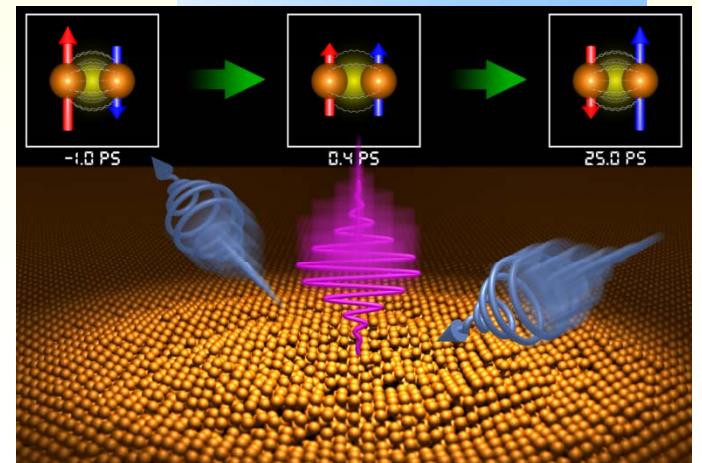
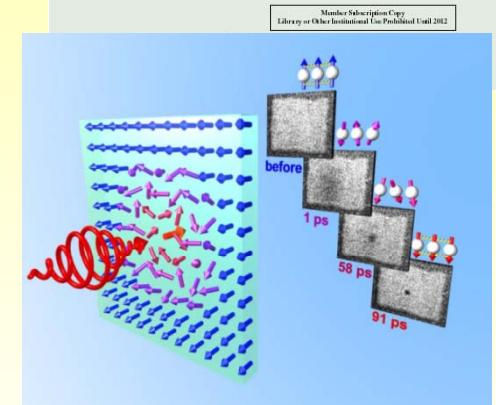
To conclude:

With light

- Coherent optical control of magnetism !
- All-optical ultrafast magnetic recording
- Novel (linear) ultrafast reversal path!
- Novel transient ferromagnetic state

Future challenges

- Femto magnetism!
- Combine chemical, magnetic spatial and time resolution: X(Z)FEL!



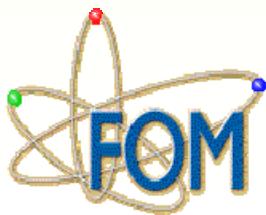
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Spectroscopy of Solids and Interfaces

Several PhD / postdoc positions available!

