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Highly sensitive electromechanical piezoresistive pressure sensors based on large-area layered PtSe₂ films

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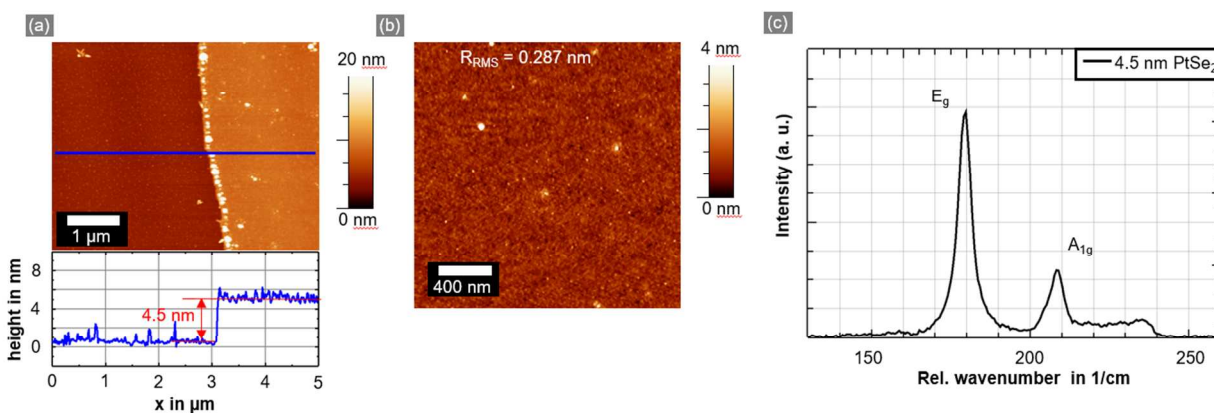
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21 KEYWORDS pressure sensors, platinum diselenide, PtSe₂, gauge factors, strain sensors
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24 ABSTRACT: Two-dimensional (2D) layered materials are ideal for micro- and
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ABSTRACT: Two-dimensional (2D) layered materials are ideal for micro- and
nanoelectromechanical systems (MEMS/NEMS) due to their ultimate thinness. Platinum
diselenide (PtSe₂), an exciting and unexplored 2D transition metal dichalcogenides (TMD)
material, is particularly interesting because its low temperature growth process is scalable and
compatible with silicon technology. Here, we report the potential of thin PtSe₂ films as
electromechanical piezoresistive sensors. All experiments have been conducted with semi-
metallic PtSe₂ films grown by thermally assisted conversion of platinum at a CMOS-compatible
temperature of 400°C. We report high negative gauge factors of up to -84.8 obtained
experimentally from PtSe₂ strain gauges in a bending cantilever beam setup. Integrated NEMS
piezoresistive pressure sensors with freestanding PMMA/PtSe₂ membranes confirm the negative
gauge factor and exhibit very high sensitivity, outperforming previously reported values by
orders of magnitude. We employ density functional theory (DFT) calculations to understand the
origin of the measured negative gauge factor. Our results suggest PtSe₂ as a very promising
candidate for future NEMS applications, including integration into CMOS production lines.

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3 Layered two-dimensional (2D) materials have extraordinary electrical, optical and mechanical
4 properties that suggest high potential for a wide range of nanoelectronics applications.^{1,2} 2D
5 transition metal dichalcogenides (TMDs) have recently been intensively investigated due to their
6 wide range of inherent electronic properties that complement graphene. Out of this family,
7 platinum diselenide (PtSe₂) is a relatively new and thus little explored TMD material. Monolayer
8 PtSe₂ is a semiconductor with a band gap of 1.2 eV,³ while bulk PtSe₂ becomes semi-metallic
9 with zero band gap.^{4,5} It has been grown epitaxially⁶ or synthesized on insulating substrates using
10 thermally assisted conversion (TAC) of predeposited platinum (Pt) films by a vaporized solid
11 selenium precursor.^{4,7} The latter process can be carried out at temperatures between 400 °C and
12 450 °C, which is compatible with standard semiconductor back end of line (BEOL) processing,⁸
13 in contrast to other 2D materials.^{9–11} We have demonstrated potential applications of TAC PtSe₂
14 in highly sensitive gas sensors and photodetectors with equal or even superior performance
15 compared to graphene and other TMD-based devices.^{7,12} Furthermore, top-down structured
16 electrical devices with PtSe₂ have been demonstrated.¹³ Here, we explore the electromechanical
17 properties of PtSe₂ in nanoelectromechanical (NEMS) pressure sensors. The devices utilize the
18 piezoresistive effect in PtSe₂, which gives rise to a change in its electrical resistance upon
19 mechanical strain. We investigate PtSe₂-membrane based pressure sensors and find sensitivities
20 that are orders of magnitude higher than those of nanomaterial-based devices,^{14–17} including
21 graphene.^{18–20} Furthermore, we obtain an average negative gauge factor (GF) of -84.8 for PtSe₂
22 from a bending beam setup. Detailed density functional theory (DFT) simulations reveal an
23 increase in the density of states (DOS) at the Fermi level of bulk PtSe₂, corroborating the
24 experimentally observed piezoresistive properties.
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3 Polycrystalline PtSe₂ films with nanometer sized grains were grown on silicon/silicon dioxide
4 (Si/SiO₂) substrates by TAC at 400 °C of predeposited 1 nm thick Pt films (Figure S1 in the
5 supporting information). Comprehensive characterization of layered PtSe₂ thin films synthesized
6 in this manner, including X-ray photoelectron spectroscopy and scanning transmission electron
7 microscopy is presented in our previous work.^{4,7} The thickness and roughness of the layers was
8 determined by atomic force microscopy (AFM, Figure 1a,b). A thickness of 4.5 nm was found
9 for the selenized 1 nm initial Pt film. Assuming an interlayer distance of 0.6 nm for PtSe₂,²¹ this
10 corresponds to 7-8 layers for the 4.5 nm thick PtSe₂ film. The root mean square (RMS)
11 roughness of the layer was 0.284 nm. The Raman spectrum of the PtSe₂ film shows two sharp
12 and prominent peaks corresponding to the E_g and A_{1g} modes of the 1T phase, attesting the
13 crystalline nature of the films (Figure 1c).⁴ After the film growth, a polymethyl metacrylate
14 (PMMA) was spin-coated on top of the PtSe₂.



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46 **Figure 1.** As-grown PtSe₂ film characterization: (a) AFM measurement of the selenized 1 nm
47 initial Pt film including height profile of the step between the Si/SiO₂ substrate to the PtSe₂ film;
48 (b) High-resolution AFM measurement displaying the RMS roughness; (c) Raman spectrum of
49 the as-grown 4.5 nm thick PtSe₂ sample.
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3 These centimeter-scale PtSe₂ thin films were transferred onto the substrate comprising of 79
4 independent pressure devices, using a typical polymer supported transfer process as described for
5 graphene.²² Each device consists of cavities covered by a patterned PtSe₂ film located in between
6 metal contacts (Figure 2a-c). The PMMA layer was not removed after the transfer in order to
7 stabilize and protect the PtSe₂ layer against environmental influences and to reduce possible
8 cross sensitivity of the sensor. Details of the fabrication processes are described in the
9 Experimental section and in the supporting information Figure S2a-d. A non-invasive Raman
10 tomography cross-section, as demonstrated recently for graphene membranes,²³ confirms the
11 presence of free-standing PtSe₂ across the cavity (Figure 2d,e) where devices with suspended
12 PtSe₂ (Figure 2e, top) and collapsed membrane (Figure 2e, bottom) were measured and can
13 clearly be distinguished with this method. The corresponding spectrum of the suspended PtSe₂
14 shows the characteristic E_g and A_{1g} modes without the Si(I) peak (Figure 2d, top), whereas the
15 spectrum taken at the bottom of the trench shows only the characteristic Si(I) peak of the Si
16 substrate (Figure 2d, bottom). This clear separation is feasible due to the 1.4 μm deep cavity,
17 which defines the distance between the suspended PtSe₂ and the Si substrate.

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38 The operation of the investigated pressure sensors can be described in the following way: the
39 PtSe₂/PMMA membranes seal the cavities and trap environmental pressure inside them.²⁴⁻²⁶ A
40 different pressure (lower or higher) outside the cavity leads to a deflection of the PtSe₂
41 membrane. This deflection induces strain in the PtSe₂ layer, resulting in a change of its
42 resistance. The chip design used in the present experiments includes devices of similar size
43 without cavities as control references. The pressure sensors, after wire-bonding and packaging
44 into a chip socket, were placed in a home-built pressure chamber where pressure can be
45 regulated from 1000 mbar (environmental pressure) down to 200 mbar (supporting information
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3 Figure S3). The pressure, humidity and temperature inside the chamber are recorded and
4 controlled with commercial sensors during the experiments. Argon (Ar) atmosphere was used to
5 exclude cross-sensitivity effects with gases and/or humidity.⁷ The current-voltage (I-V)
6 characteristics of the PtSe₂ device for environmental pressure shows good linearity, i.e. the
7 transferred PtSe₂ forms good Ohmic contacts with metal (gold) electrodes (supporting
8 information Figure S4). The measured data is presented in Figure 2f along with the
9 simultaneously measured pressure (dashed red line, right y-axis). The resistance across the PtSe₂
10 sensors decreases with applied increasing strain on the membrane (caused by the decrease in
11 pressure, supporting information Figure S5a). A pressure difference of 800 mbar between the
12 outside and the inside of the cavity leads roughly to 7% of change in resistance.
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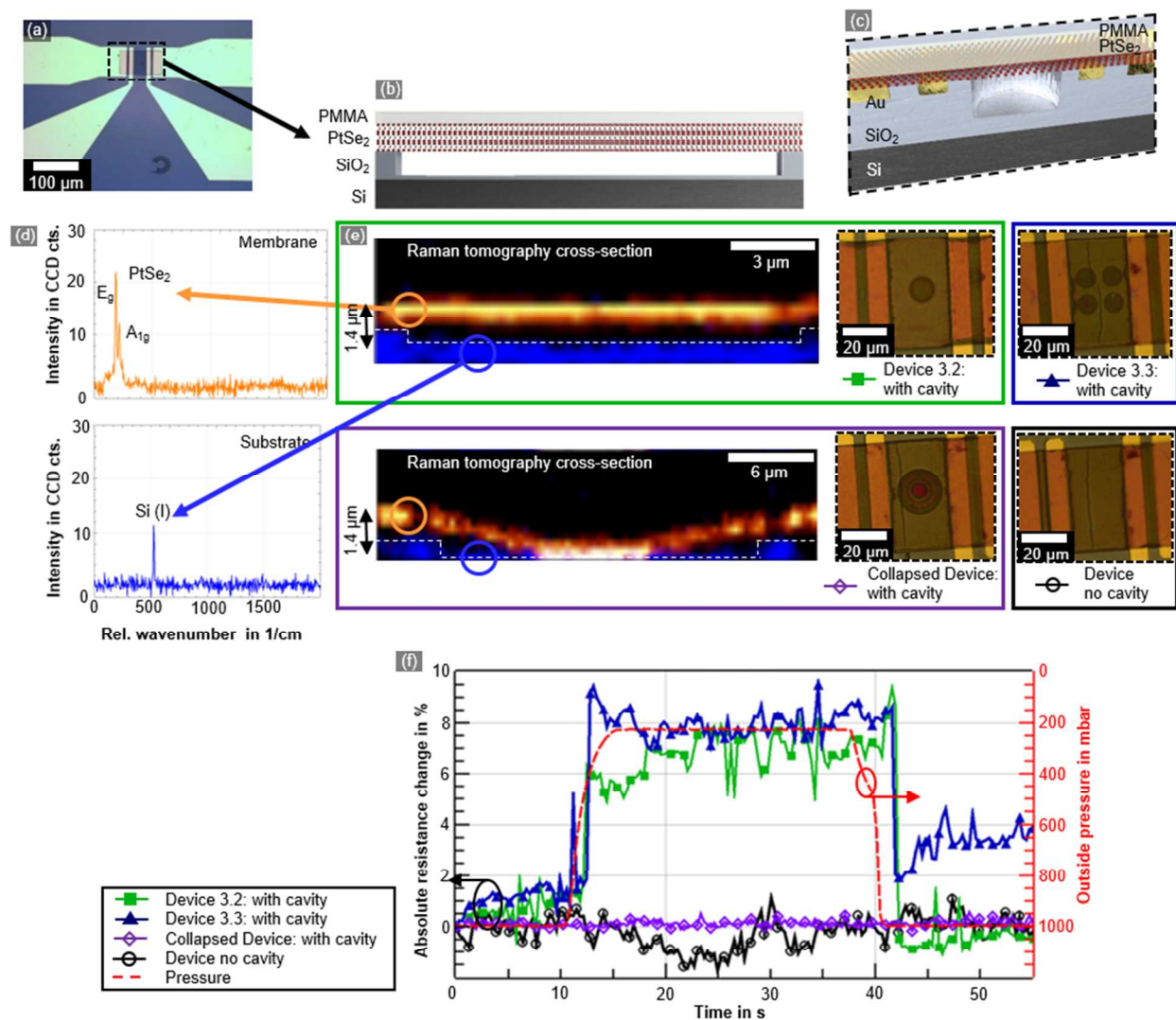


Figure 2. Pressure sensor fabrication and measurement. (a) Optical micrograph of one device with Au contacts and the PtSe₂ channel across the cavity area. (b) Schematic cross-section of the cavity area with suspended PtSe₂. (c) Schematic of one device including contacts and cavity area. (d) Raman point spectrum of the PtSe₂ membrane area and the Si substrate. (e) Raman tomography cross-sections of one device with suspended PtSe₂ (device 3.2) and one with collapsed membrane (Collapsed Device). On the right are the corresponding optical micrographs of the cavity areas (device 3.2 and the collapsed device) as well as one other working device (device 3.3) and one device without cavity (device no cavity) as reference. (f) Resistance change

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3 in % for one pressure cycle (vented, pumped down to 200 mbar, vented) of the four devices
4 shown in (e) (left y-axis) and pressure on the outside of the membrane (right y-axis).
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8 The results of four devices are displayed in Figure 2f: two intact membranes on cavities (filled
9 green squares and filled blue triangles), one with a collapsed membrane (empty purple diamond)
10 and one device without cavity for reference (empty black circles). A slight decrease (increase) in
11 humidity during evacuation (venting) of the chamber is commonly observed for all
12 measurements (supporting information Figure S5b), but does not give rise to changes in the
13 resistance. Cross-sensitivity towards humidity can thus be excluded from having any substantial
14 influence on the device behavior. We conclude that the observed resistance change is primarily
15 caused by pressure (i.e. strain) changes during the measurements.
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27 The absolute sensitivity (S) of a piezoresistive device is calculated using equation (1), where
28 $\Delta R/R$ is the absolute value of the relative resistance change and ΔP is the pressure change during
29 the sensor operation.¹⁸
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$$S = \frac{\Delta R}{R \cdot \Delta P} \quad (1)$$

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39 A total of 19 working devices were measured on four different chips through 69 different
40 measurements (Figure 3a). A sensitivity of $1.05 \times 10^{-4} \text{ mbar}^{-1}$ was extracted from Figure 2f. An
41 average of $5.51 \times 10^{-4} \text{ mbar}^{-1}$ was calculated for all measured devices (best device not included).
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46 The results were compared to some other NEMS pressure sensors which operate in a similar
47 pressure range (Figure 3b,c). These use gallium arsenide (GaAs),²⁷ Si nanowires (Si NWs),^{15,16}
48 single- and multi-wall carbon nanotubes (SWNT and MWNT)^{14,17} and graphene^{18–20} as the
49 piezoresistive materials. However, for most of these (except Smith *et al.*¹⁸ and Fung *et al.*¹⁷) a
50 separate membrane is needed which is treated, or where the piezoresistive materials are grown or
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3 applied, in a separate fabrication step. This adds additional complexity to the fabrication process,
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5 which is not the case for the PtSe₂-based pressure sensors, because the material acts as
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7 membrane and sensor at the same time. The thickness of other material is limited due to their
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9 mechanical stability. The measured sensitivity outperforms the best reported 2D and low
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11 dimensional materials based pressure sensors by a factor between 3 and 70 (Figure 3b). The
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13 SWNT-based pressure sensor¹⁴ shows an equal sensitivity, however the membrane areas are
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15 nearly two to four orders of magnitude larger than the PtSe₂ membrane. To take this into account
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17 the sensitivity can be normalized to the membrane area to receive a direct comparison between
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19 the devices independent of their membrane size.^{28,18} Here, the PtSe₂-based pressure sensor shows
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21 nearly 2 to 5 orders of magnitude higher sensitivity than the other devices (Figure 3c, see also
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23 supporting information Table S6).
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29 We further point out that one of the devices showed an exceptionally high sensitivity of $1.64 \times$
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31 $10^{-3} \text{ mbar}^{-1}$ (Figure S5c). Even though it could not be reproduced in other devices, this value was
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33 repeatedly measured for that particular device (Figure 3a, device 2.1). The sensitivity and
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35 normalized sensitivity of this device are plotted in Figure 3b,c (black star). When normalized to
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37 the membrane area, this device is over 3 orders of magnitude more sensitive than the highest
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39 performing devices shown in Figure 3c (a SWNT¹⁴ and a graphene membrane¹⁸ based device),
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41 and 4 to 6 orders of magnitude more sensitive than other reported devices.
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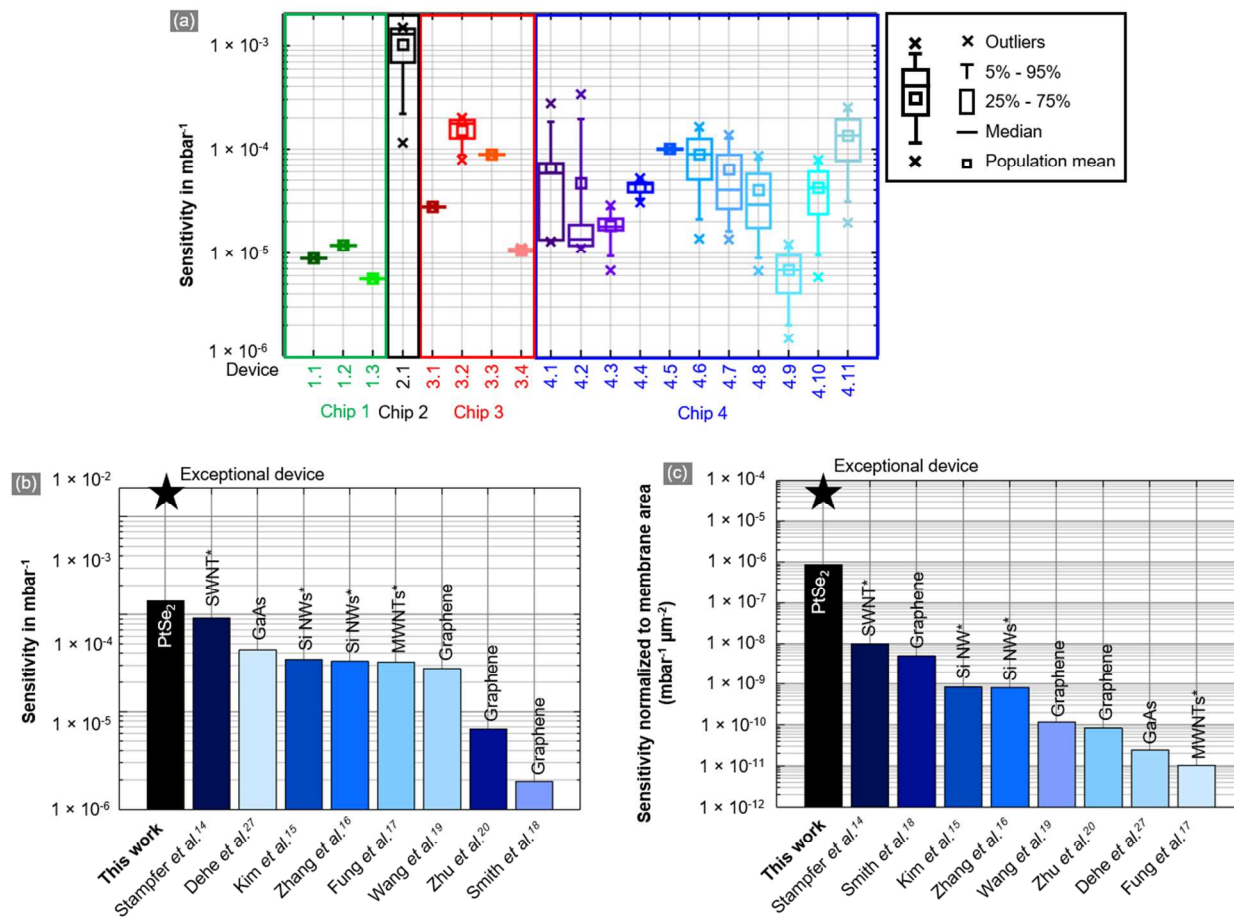


Figure 3. Comparison of all measured chips and devices as well as comparison of various pressure sensor devices using different materials for piezoresistive detection. (a) Box and whiskers plot of all measured chips, devices and repeated measurements; (b) Comparison of the sensitivity in mbar^{-1} with other pressure sensors; (c) comparison of the sensitivity normalized by membrane area in $\text{mbar}^{-1} \mu\text{m}^{-2}$ with other pressure sensors. In plots (b) and (c), devices which require a separate membrane are indicated with a * and the device with exceptionally high sensitivity is indicated by a black star.

A bending beam setup was used to measure the piezoresistive gauge factor of the 4.5 nm and 9 nm thick PtS₂ films, which consisted of 7 to 8 and approximately 15 layers, respectively. The

bending beam setup enables measurements of well-defined strain fields in a very simple way (for details see supporting information). The gauge factor (GF) can be extracted using equation (2).²⁹

$$GF = \frac{\Delta R}{R_0 \cdot \varepsilon} \quad (2)$$

Here, ε is the strain in the bending beam at the position of the sensor ($\varepsilon = 0.04\%$). A negative gauge factor (GF) of -84.8 was obtained (Figure 4a, supporting information section S7 and S9). The measured resistance change is shown in Figure 4b and in the supporting information Figure S9b for two PtSe₂ layer thicknesses and masses of 2 kg and 0.5 kg (Figure S9c in the supporting information). A measurement of multiple strain cycles with a mass of 2 kg is shown in Figure S9d in the supporting information. In some measurements, the resistance did not immediately return completely to its initial values after removing the mass (Figure 4b, 4.5 nm PtSe₂ device, Figure S9b, d in the supporting information). This may be attributed to slip of the device or the electrical contacts during measurements due to non-optimal adhesion to the beam. Furthermore, a slight overshoot followed by oscillations was observed for some devices right after applying the mass. These may be interpreted as natural oscillations due to the manual application of the mass on the beam. Data from a commercially available metal strain gauge is used as a reference (Figure 4b). A decrease in resistance with increasing strain on the PtSe₂ film corroborates the integrated pressure sensor measurements (supporting information Figure S9a). The GF of the PtSe₂ films exceeds the GF of the metal strain gauge (GF of 2)²⁹ and of graphene (GF of 2-6).^{24,30,31} While similar or higher GF have been found for doped silicon (GF of -100 to +200)²⁹ or mono- and trilayer MoS₂ (GF of -148 and -43.5).³²

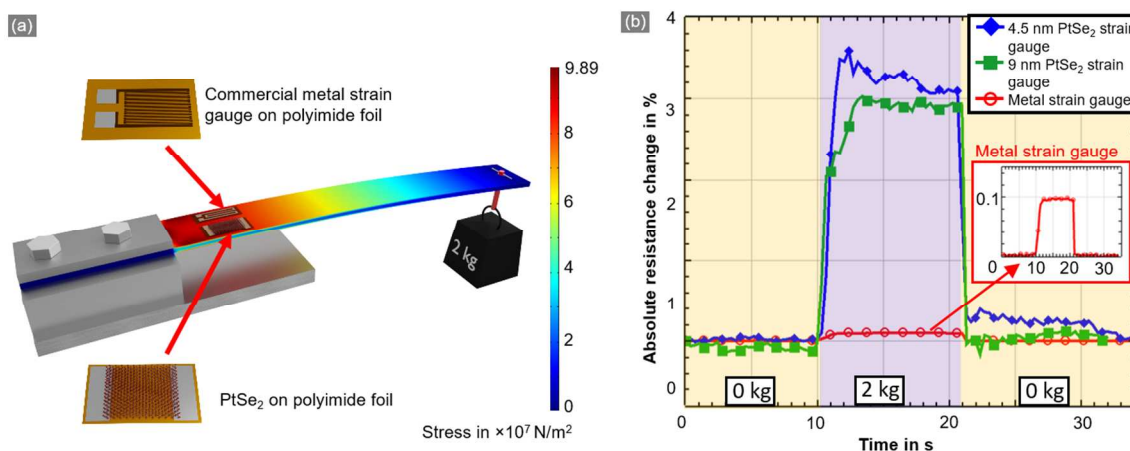


Figure 4. (a) Bending beam setup (cantilever beam) with applied PtSe₂ and commercially available metal strain gauges including stress simulation with the applied weight; (b) Electrical readout signal during the measurement with an absolute resistance change against time showing two PtSe₂ devices (of 4.5 nm PtSe₂, diamonds filled in blue and 9 nm PtSe₂, squares filled in green) and a metal strain gauge (empty red circles) as reference with an enlarged view in the inset.

DFT calculations were conducted to gain insights into the observed resistance decrease with increased strain in PtSe₂ films. The model was designed to represent the membrane covering the cavities of the PtSe₂ pressure sensors. Here, the deformation can be thought of as bi-axial strain acting on the PtSe₂ membrane. The model was fully optimized in terms of the lattice vectors and atomic positions. In the bulk system, which is similar to the experimental conditions in this work, biaxial strain means that the layers are stretched or compressed in plane. As bulk PtSe₂ is metallic, the density of states (DOS) at the Fermi level is found to increase under biaxial tensile strain (ϵ_a) and reduces under compression, as shown in the supporting information Figure S10a. This result is in accordance with the experimental findings, where the resistivity decreases with increased ϵ_a .

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3 Next, we have considered the effect of the interlayer strain and compression (ϵ_c) on the DOS.
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5 The results are shown in the supporting information Figure S10b. The compression of the
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7 interlayer distance has the opposite influence on the system and it slightly reduces the DOS at the
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9 Fermi level. Finally, we have considered the two effects acting in tandem in the following way:
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11 for in-plane stretching, there is compression out of the plane, and vice versa (Figure 5a). The
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13 results for the combined effect are shown in Figure 5b. Once again, the DOS at the Fermi level is
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15 increased for the in plane stretching and compression in the out-of-plane direction. However, the
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17 effect is reduced compared to just ϵ_a , due to the opposite effect of the interlayer compression, as
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19 shown in the supporting information Figure S10b. In comparison to the significant increase of
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21 the DOS at the Fermi level upon strain the band structure shows only minor changes (Figure
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23 S10c in the supporting information) not indicating strong differences in the effective masses
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25 under strain.
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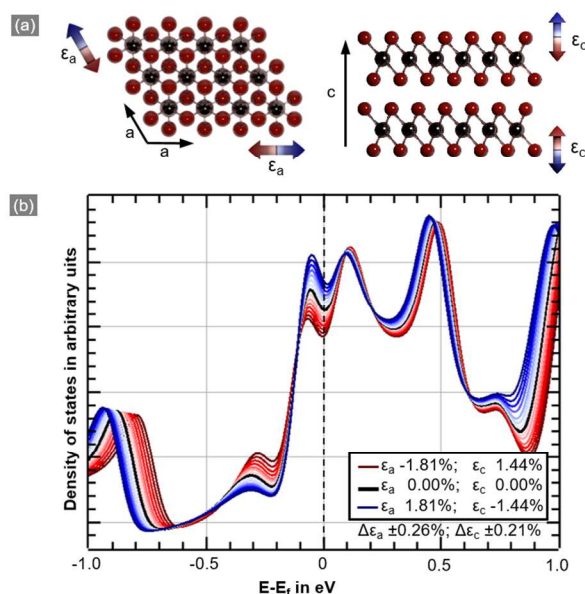


Figure 5. (a) Top and side view of the PtSe₂ bulk shown with the in-plane, a and out-of-plane, c lattice vectors, and the direction of the applied strain (ϵ_a and ϵ_c). (b) Density of states close to the Fermi level (shifted to zero) under applied tensile strain and compression.

An increase in DOS has also been reported for other single and few layer semiconducting TMDs, like MoS₂,³³ MoSe₂,³⁴ WSe₂³⁴ and also PtSe₂.³⁵ The semiconducting materials exhibit a decrease in band gap and even a transition to metallic behavior with applied strain, resulting in a negative gauge factor. For more detailed future studies the influence of many other parameters should be considered and investigated, including the adhesive interface between the steel beam and the polyimide foil, the thickness and mechanical properties of the foil and the PMMA layer covering the PtSe₂. Also, the crystallographic orientations, thickness and other material properties of the TAC-grown PtSe₂ film need to be investigated and considered in the simulations. Nevertheless, the results of the simulations show a clear indication of an increase in DOS, leading to a decrease in resistance under the applied strain, which is the most plausible cause of the experimentally extracted negative GF.

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3 We have demonstrated PtSe₂-based piezoresistive pressure sensors with high gauge factors and
4 very high sensitivity. The experiments were conducted on large-area PtSe₂ thin films grown by a
5 scalable method of thermal conversion of metallic films at a low temperature of 400°C. A
6 negative gauge factor of approximately -85 was obtained from a PtSe₂ strain gauge through
7 bending beam experiments and can be understood in terms of a strain-induced increase in the
8 DOS, as verified through DFT simulations. Furthermore, integrated piezoresistive pressure
9 sensors were fabricated using PtSe₂/PMMA membranes. The PtSe₂ pressure sensors showed very
10 high sensitivity, outperforming piezoresistive pressure sensors based on other materials by at
11 least a factor of 82. When normalized to the active membrane area, the sensitivity of our devices
12 is almost two to five orders of magnitude larger than values reported for any other technology to
13 date. One exceptional device has an even higher sensitivity of $1.64 \times 10^{-3} \text{ mbar}^{-1}$ and outperforms
14 existing technologies by an even greater margin. These results suggest that layered 2D PtSe₂ has
15 great potential for future piezoresistive NEMS applications, beyond the demonstrated pressure
16 sensors. These devices have a footprint that is up to four orders of magnitude smaller than other
17 pressure sensors while showing a similar or even higher sensitivity. Further down-scaling to
18 smaller cavities will improve the stability and yield of the membranes, and allow decreasing the
19 membrane thickness to maintain the high sensitivity. The low growth temperature and general
20 CMOS material compatibility makes PtSe₂ highly attractive for future silicon technology
21 integration.
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Methods

Materials Synthesis and analysis

Pt layers with a thickness of 1 and 2 nm were sputter deposited on $1.5 \times 1.5 \text{ cm}^2$ Si/SiO₂ substrates using a Gatan coating system (Gatan 682 PECS) with a 19 mm diameter Pt target. A TAC process was used to selenize the pre-deposited Pt layers, as described in detail in references.^{4,7} The cross-sectional Raman measurements were conducted using a WITec alpha 300R system with a 532 nm laser as source of excitation. A laser power of 75 μW was used. For obtaining single Raman spectra an 1800 g/mm grating was used, while the depth scan (cross-section) was conducted with a 300 g/mm grating. The system is limited to a lateral and vertical resolution of 300 nm and 900 nm, respectively.

Strain gauge fabrication

The strain gauges were fabricated by transferring as-grown PtSe₂ films onto pre-patterned polyimide foil substrates. A typical polymer-supported transfer method was used for the film transfer. 200 nm thick PMMA (Allresist AR-P 649.04) was spun on the as-grown PtSe₂ samples at a speed of 2000 rpm for 50 s. After spin-coating, the samples were annealed for 5 min at 80°C to cure the PMMA. 2 molar KOH was used to etch away the SiO₂ layer and release the PMMA-covered PtSe₂ films from the substrate. The PMMA-covered PtSe₂ films were cleaned in DI-water, followed by a transfer to the pre-patterned (500 nm thick copper contacts from evaporation) polyimide foil. After drying, the foil was glued (super glue Z70 by HBM) onto the steel beam (size of $300 \times 30 \text{ mm}^2$, 3 mm thick), 200 mm away from the loading point. A commercially available metal strain gauge was glued next to the PtSe₂ strain gauge. Cables were then soldered onto the contact pads for electrical connections.

Pressure sensor fabrication

A Si chip ($7 \times 7 \text{ mm}^2$) with $1.6 \text{ }\mu\text{m}$ thick SiO_2 was used as a device substrate (Figure S2a in the supporting information). The cavities have circular geometry (radii between $2 \text{ }\mu\text{m}$ and $20 \text{ }\mu\text{m}$) with arrays of 1 to 15 cavities and rectangular geometries (length of $20 \text{ }\mu\text{m}$ and width between $2 \text{ }\mu\text{m}$ and $4 \text{ }\mu\text{m}$) in arrays between 1 and 3 cavities. They were structured by photolithography, followed by Ar- and CHF_3 -based reactive ion etching (RIE, 200 mW, 40 mTorr) to create a vertical profile with $1.4 \text{ }\mu\text{m}$ of depth (Figure S2b in the supporting information). The contact regions were defined by photolithography, followed by a RIE process to embed the contacts and self-aligned metal evaporation (Cr/Au, 40 nm/400 nm, Figure S2c in the supporting information). The PtSe_2 films grown from 1 nm initial Pt films were transferred onto the device substrate using a PMMA-supported transport method as described above (Figure S2d in the supporting information). Lastly, the chip was wire bonded into a 44-pin PLCC chip package using $25 \text{ }\mu\text{m}$ diameter Au wire.

Simulation

All calculations were performed using the DFT method as implemented in the VASP plane wave code.^{36,37} Ion-electron interactions were described by projector augmented wave approximation.³⁸ We have employed Perdew–Burke–Ernzerhof (PBE)³⁹ functional under the generalized gradient approximation. Van der Waals interaction corrections were included using the D3 approach proposed by Grimme.⁴⁰ Cutoff energy was set to 600 eV and the convergence threshold for residual force was $0.01 \text{ eV }\text{\AA}^{-1}$. Brillouin zone integration was carried out at $6 \times 6 \times 6$

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3 Monkhorst–Pack k-grids. Spin-orbit coupling (SOC) was included in the electronic structure
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5 calculations. Bulk PtSe₂ was fully optimized (lattice vectors and atomic positions) resulting in
6
7 the following lattice vectors: $a = b = 3.788 \text{ \AA}$ and $c = 4.790 \text{ \AA}$, and a distance between Se atoms
8
9 in the Se-Pt-Se sandwich of 2.550 \AA . This is in good agreement with the experimental data by
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11 Wang et al.,⁶ with $a = b = 3.70 \text{ \AA}$ and the distance between Se atoms of 2.53 \AA .
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ASSOCIATED CONTENT

Supporting Information

Schematic of the growth of PtSe₂ with additional Raman spectra. The fabrication process is described in a schematic and the measurement chamber is shown on a schematic. I-V characteristic of the unstrained pressure device is shown as well as some additional electrical characterization data of the pressure devices. A table of the compared pressure sensors from literature is shown. The strain gauge measurements are described in detail, including the I-V characteristic of the device as well as some further electrical measured data. Further DFT calculations are shown for bulk and for three layer PtSe₂. (PDF)

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Author Contributions

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20 21 ABBREVIATIONS

22
23 2D, two dimensional; PtSe₂, Platinum diselenide; TMD, transition metal dichalcogenides;
24 MEMS, microelectromechanical systems; NEMS, nanoelectromechanical systems; DFT, density
25 function theory; TAC, thermally assisted conversion; Pt, Platinum; GF, gauge factor; DOS,
26 density of states; Si, silicon; SiO₂, silicon oxide; AFM, atomic force microscopy; RMS, root
27 mean square; PMMA, polymethyl metacrylate; Ar, argon; I-V, current-voltage; GaAs, gallium
28 arsenide; Si NW, silicon nanowire; SWNT, single wall nanotube; MWNT, multi wall nanotube.
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