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Morphology Tunable Synthesis of Intrinsic Room-Temperature Ferromagnetic γ -Fe₂O₃ Nanoflakes

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ABSTRACT: Intrinsic two-dimensional (2D) ferromagnets with room temperature ferromagnetism and air stability are highly desirable for spintronic applications. Unfortunately, the experimentally observed intrinsic 2D ferromagnetic materials are very less, and there are fewer materials that have both room temperature ferromagnetism and air stability. Here, intrinsic room-temperature ferromagnetic γ -Fe₂O₃ nanoflakes with different shapes have been successfully grown by the chemical vapor deposition method. Firstly, triangular γ -Fe₂O₃ flakes epitaxially grow on the α -Al₂O₃ surface at low temperature, and the flakes grow gradually in

size and finally form a large-scale γ -Fe₂O₃ film as increasing the growth time. This growth process belongs to thermodynamic control. However, the triangular γ -Fe₂O₃ flakes sequentially transform into stellated, petaloid, and dendritic crystalloids when enhancing the temperature of precursor powder. This transition reveals the growth process from thermodynamically to kinetically dominated control and successfully realizes the morphology tunable synthesis recipe. Additionally, these γ -Fe₂O₃ flakes are ferromagnetic in the range of 5 K to 300 K and own out-of-plane easy magnetization axis. Moreover, compared with the triangular and stellated γ -Fe₂O₃ flakes, the petaloid and dendritic γ -Fe₂O₃ flakes show enhanced coercive force. The γ -Fe₂O₃ nanoflakes with diverse shapes synthesized by controlling temperature have similar resistivity. The synthesis of γ -Fe₂O₃ flakes with different shapes and properties paves the way for the preparation and application of next generation electronic and spintronic devices.

Keywords: γ -Fe₂O₃ Nanoflakes; Morphology tunable synthesis; Room-temperature ferromagnetism; Coercive force; Resistivity

Introduction

Ferromagnets (FMs) advance greatly the development of state-of-the-art electronic and spintronic devices.¹⁻³ The new generation of memory devices triggers ever-increasing requirements on the dimension of FMs. Recent appearance of two-dimensional (2D) layered materials provides the feasibility of small size and high integration for novel memory devices.⁴ ⁵ The advantage of these 2D layered materials rests with their in-plane atoms linked by strong chemical bonds and their out-of-plane layers linked by a weak van der Waals (vdWs) force.^{6,7} Such an extremely low dimension not only induces unprecedentedly abundant physics but also

reveals a magnificent innovation in ever-smaller electronic and spintronic devices.^{8,9} However, such an intriguing 2D magnetism remains an elusive dream in practice, and these stable 2D crystals are not triggered in real yield yet.¹⁰

Diverse 2D layered magnetic materials have currently been reported experimentally, e.g. CrI₃,¹¹⁻¹⁶ Fe₃GeTe₂,^{1, 17-19} and Cr₂Ge₂Te₆,²⁰, and accordingly, these 2D magnetic materials have been widely used in electronic and spintronic devices due to their exotic property. Nevertheless, the most studied 2D layered FMs so far are unstable in atmosphere, and their Curie temperature (T_C) is far below room temperature.²¹ Recently, a few 2D non-layered materials were achieved using the chemical vapor deposition (CVD), which supposedly solved the issue of air stability. For instance, Cr₂S₃ is an air-stable ferrimagnetic semiconductor.²²⁻²⁴ In particular, CrSe²⁵ and Cr₂Te₃²⁶ behave close to room-temperature T_C (~280 K) in experiments. These experimental studies are of particular interest in that direct synthesis of room-temperature FMs through the CVD method is key for promoting device utility.²⁷ Despite that many kinds of 2D layered and non-layered FM-based devices have been demonstrated at low or high temperatures, they are not yet adequately warm to room temperature for electronic and spintronic device applications. A 2D FMs possessing simultaneously ultrathin nature, air stability, and high T_C is thus timely and ultimately important for a new generation of electronic and spintronic devices.

As an iron oxide, 2D non-layered ϵ -Fe₂O₃ represents such a kind of FM materials in view of its ultrathin nature and room-temperature T_C .²⁸ Iron oxide is a species of common material rich in natural resources and nontoxicity,²⁹ which involves various phases, such as α , β , γ , and ϵ .³⁰ Among these phases, γ -Fe₂O₃ is the second most common Fe₂O₃ polymorph in nature that could be applied in tumor treatment,³¹ lithium-ion battery,³² recoverable catalyst,³³ gas sensor,³⁴

and magnetic study.^{35,36} It has a cubic crystal structure of an inverse spinel-type and crystallizes in the $P4_132$ space group with $a = 8.351 \text{ \AA}$ and consists of one-third of the tetrahedral sites and two-thirds of the octahedral sites bearing Fe^{3+} cation. Importantly, $\gamma\text{-Fe}_2\text{O}_3$ has a spinel structure with two magnetic sublattices and exhibits ferromagnetism,^{30,37} making it a preferred oxide to fulfill air stability and room-temperature ferromagnetism for ultrathin electronic and spintronic device applications. However, the synthesis of ultrathin room-temperature ferromagnetic $\gamma\text{-Fe}_2\text{O}_3$ materials is scarcely reported due to the difficulty in controlling the growth process and the easy mutual conversion between $\gamma\text{-Fe}_2\text{O}_3$ and $\alpha\text{-Fe}_2\text{O}_3$ or oxyhydroxide.^{29,38}

Here, we report the successful synthesis of ultrathin $\gamma\text{-Fe}_2\text{O}_3$ nanoflakes using a facile and efficient growth strategy based upon the CVD method,^{2,39} and demonstrate that ultrathin $\gamma\text{-Fe}_2\text{O}_3$ nanoflakes exhibit diverse morphology as the growth temperature is tuned. Importantly, all these nanoflakes with different shapes show room temperature ferromagnetism. At low temperature, triangular $\gamma\text{-Fe}_2\text{O}_3$ flakes appear and the flakes grow gradually in size and finally form a large-scale $\gamma\text{-Fe}_2\text{O}_3$ film as increasing the growth time. This growth process belongs to thermodynamic control. The triangular $\gamma\text{-Fe}_2\text{O}_3$ flakes sequentially transform into stellated, petaloid, and dendritic crystalloids as increasing the temperature of precursor powder. This transition reveals the growth process from thermodynamically to kinetically dominated control. Moreover, compared with the triangular and stellated $\gamma\text{-Fe}_2\text{O}_3$ flakes, the petaloid and dendritic $\gamma\text{-Fe}_2\text{O}_3$ flakes show enhanced coercive force. The $\gamma\text{-Fe}_2\text{O}_3$ nanoflakes with diverse shapes synthesized by controlling temperature have similar resistivity. The intriguing finding of ultrathin, air-stable, and multifunctional room-temperature 2D non-layered $\gamma\text{-Fe}_2\text{O}_3$ FM and its growth recipe opens up a novel avenue in the practical application of next-generation electronic,

spintronic, and memory devices.

Results and Discussion

γ -Fe₂O₃ can be obtained directly from the oxidation of iron, which can be achieved using a simple CVD system (Figure S1a). Iron and cobalt mixed powder was loaded in a corundum boat, which was placed in the center of the furnace temperature zone. Cobalt can be used as a fluxing medium or catalyst in the synthesis process. The substrate [SiO₂/Si, mica, or sapphire (α -Al₂O₃)] was placed face down on the corundum boat; sapphire was more suitable for γ -Fe₂O₃ nanoflake growth (Figures S1b–d). The nanoflakes were then characterized by optical microscopy (OM), scanning electron microscopy (SEM), Raman spectroscopy, and X-ray photoelectron spectroscopy (XPS). As shown in Figure 1a, the triangular γ -Fe₂O₃ nanoflakes grown on sapphire exhibited a uniform orientation, which might be related to the substrate. To verify this point of view, statistics was performed. Figure S2a shows the photograph of a 2-inch α -Al₂O₃ substrate with a unique polished surface [c-plane (0001)] and a cutting reference edge [α -Al₂O₃ (11 $\bar{2}$ 0)]. The SEM images of the triangular γ -Fe₂O₃ nanoflakes (upper) and the cutting reference edge (lower) are illustrated in Figure S2a, and the green and black arrows represent the crystal edge (zigzag direction) of a γ -Fe₂O₃ nanoflake and the cutting reference edge [α -Al₂O₃ (11 $\bar{2}$ 0)], respectively. θ is the angle between the crystal edge of γ -Fe₂O₃ nanoflakes and the cutting reference edge. The atomic models of γ -Fe₂O₃ and α -Al₂O₃ simulated this growth stacking mode (Figure S2b). The orientations of γ -Fe₂O₃ nanoflakes were strongly coupled with single-crystal facets (Figure S2c), indicating that the γ -Fe₂O₃ nanoflakes epitaxially grew on the α -Al₂O₃ surface, which was attributed to their similar lattice parameters

and lowest formation energy for γ -Fe₂O₃ nanoflakes at $\theta \approx 0$. Our paper demonstrates that the epitaxial growth recipe could promote uniform orientations of γ -Fe₂O₃ grown on single-crystal sapphire. This work is similar to some works that have determined mono-orientated 2D materials by using low-crystalline substrate facets.^{40, 41} The triangular nanoflakes were then transferred on a copper grid (Figure 1b) to investigate the structure of γ -Fe₂O₃. As revealed in the high-resolution transmission electron microscopy (HRTEM) image (Figure 1c) collected at the square location in Figure 1b, atoms were distributed in accordance with the honeycomb lattice, and the (0 $\bar{2}2$) interplanar spacing was ~ 0.3 nm, in fair agreement with the top view of the γ -Fe₂O₃ structure along [111] axis. Given that transmission electron microscopy (TEM) and HRTEM images were collected in the same direction, the edge of the triangular γ -Fe₂O₃ nanoflakes corresponded to the zigzag direction of the crystal, as marked by the white arrows in Figures 1b and c. Figure 1d shows the corresponding fast Fourier transform (FFT) pattern of the HRTEM image, which displays one set of sixfold symmetry diffraction spots. These results agree with those of previous works,⁴²⁻⁴⁴ proving that the triangular nanoflakes are definitely γ -Fe₂O₃ crystals. The Raman spectrum (Figure 1e) of the as-grown triangular nanoflakes exhibits three peaks at ~ 356 , 512, and 650–720 cm⁻¹, except for the contribution from the substrate, they were recognized to arise from T_{2g} , E_g , and A_{1g} , respectively. Two broad peaks appeared at high wavenumbers (~ 1200 and 1400 cm⁻¹), which might be contributed by the magnon mode.^{38, 45, 46} These characteristic peaks are highly consistent with reports on γ -Fe₂O₃ crystals.⁴⁷⁻⁵² XPS measurements were also performed to explore the chemical state of Fe and O elements. The high-resolution region of Fe 2p (Figure 1f) can be divided into two main peaks at binding energies of 710.71 (Fe 2p_{3/2}) and 724.25 eV (Fe 2p_{1/2}) together with two shake-

up satellites at 717.96 and 731.85 eV, which are consistent with the typical values for Fe (III).⁵³⁻

⁵⁵ All the peaks were fitted using the Gauss–Lorentz (ratio: 40%) combination formula. The deconvolution spectrum of O 1s (Figure 1g) shows a peak at 531.07 eV, which can be fitted using the Gauss–Lorentz (ratio: 30%) combination formula. The OM, SEM, Raman, and XPS results strongly proved the synthesis of triangular γ -Fe₂O₃ nanoflakes.

A series of morphology images for samples in different growth stages was collected to explore the synthesis process of triangular γ -Fe₂O₃ nanoflakes. Figures 2a–g show the SEM images of γ -Fe₂O₃ nanoflakes synthesized with different growth times. In the initial stage, the γ -Fe₂O₃ nanoflakes showed three related dendrites along specific diffusion pathways (Figure 2a). We inferred that these dendrites were related to the initial clusters. Then, the crystal grew along the edges of the three dendrites. After 5 min of growth, the edges of the crystal displayed a sawtooth shape (Figure 2b). Subsequently, the γ -Fe₂O₃ nanoflakes further grew into a nearly complete triangular shape (Figures 2c and d). As shown in Figures 2e–g, the triangular γ -Fe₂O₃ nanoflakes rapidly grew up and gradually spliced together, eventually forming a large-scale film, this growth process belongs to thermodynamic control.^{56, 57} The AFM contour images confirmed this phenomenon. Figure S3a shows the AFM image of the γ -Fe₂O₃ nanoflakes at the initial stage; three branches were growing in three directions, which were formed with tiny triangular sheets. As shown in Figure S3b, the γ -Fe₂O₃ nanoflake obtained through a longer growth time had a triangular shape with three jagged edges. The surface of the γ -Fe₂O₃ nanoflake was uneven, indicating that it was in an immature growth stage. Figure S3c depicts the AFM image of a triangular nanoflake with three sharp edges, in which the traces of three growth pathways (as mentioned in the initial stage) were still faintly visible on its surface. The

thickness of the nanoflakes did not change dramatically during the growth process. Figure 2h shows the relationship plot between nanoflake size and growth time; each data point displayed with a standard error of the mean was counted from multiple nanoflakes. Comparison of the slopes of the two pieces of data (marked with red and blue arrows) indicated that the crystal size increased slowly in the early stage. However, the growth rate was considerably accelerated when the growth time exceeded 11 min. On the basis of the above results, the schematic of the growth recipe of triangular γ -Fe₂O₃ nanoflakes is illustrated in Figure 2i. Many dangling bonds were exposed on the edge of the dendrites at the initial growth stage. Subsequently, iron and oxygen atoms were preferentially navigating to a stable state or an energetically more favorable location along the edges or corners of the dendrites.⁵⁸ The fluxing medium (Co) may play an active role in crystal growth.⁵⁹ Accordingly, the crystal growth laterally from the dendrite edge would be promoted, facilitating the synthesis of regularly shaped γ -Fe₂O₃ nanoflakes. As the growth time extended, the nanoflakes grew gradually in dimension and merged, forming a large-scale γ -Fe₂O₃ film. These findings are similar to previous reports on other 2D materials.^{58,}

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Synthesis of 2D materials by using the CVD method involves many inconstant growth parameters, such as substrate, transport gas, growth time, and cooling rate. In this work, suitable substrate and transport gas (Ar) were selected, and the influence of growth time on triangular nanoflakes was clarified. Furthermore, the precursor temperature acts as a dominant factor for the nanoflake shape and size in the crystallization and growth stage.⁶⁵ Figure 3 shows the OM and SEM images of the γ -Fe₂O₃ nanoflakes grown at different temperatures. The overall results indicated that the growth of γ -Fe₂O₃ nanoflakes followed several novel growth

models. Figures 3a–d show some OM images of the four nanoflake models in the initial stage (growth time: ~4 min). The growth temperatures were 930 °C, 980 °C, 1030 °C, and 1080 °C. With prolonged growth time (~15 min), γ -Fe₂O₃ nanoflakes could grow into mature shapes: triangular, stellated, petaloid, and dendritic shapes (Figures 3e–h). As described in the previous section, the triangular γ -Fe₂O₃ domains diverged out along the three diffusion pathways. Nonetheless, the γ -Fe₂O₃ domains tended to grow two sets of related dendrites (six diffusion pathways) derived from the same nucleated cluster when the temperature was set to 980 °C. Finally, the domains grew into stellated crystals, including two triangular crystals with a specific angle of 180°, which were synthesized on the basis of the same nucleated cluster. The γ -Fe₂O₃ domains changed into a petaloid shape when the temperature was increased to 1030 °C. Figures 3c and g present this novel diffusion-limited growth model. The number of dendrites derived from the same nucleated cluster increased significantly, and the nanoflakes grew into petaloid ones along more than sixfold diffusion pathways (displaying a radial pattern). Figures 3d and h represent an exceptional growth model, which was dominated by dendrite growth when the temperature was increased to 1080 °C. In the initial stage (Figure 3d), the γ -Fe₂O₃ nanoflakes grew into a fan-shaped blade. As the growth time increased, the crystal grew chaotically along the dendrites (Figure 3h). The high-magnification OM (Figures 3i–l) and SEM (Figures 3m–p) images of these nanoflakes reveal the shape and size of these triangular, stellated, petaloid, and dendritic nanoflakes. In this work, we have achieved controlled growth of nanoflakes in various shapes, and the temperature is a pivotal factor for tunable morphology.⁶⁵ The growth details will be analyzed in the following paragraphs.

To explore the growth path of these shaped γ -Fe₂O₃ nanoflakes, their surface AFM contour

images were collected. Figures 4a–d show the AFM images of triangular, stellated, petaloid, and dendritic nanoflakes with ultrathin thicknesses of ~4.61, 4.15, 8.67, and 5.31 nm, respectively. The high-magnification AFM images (Figures 4e–h) display the surface details of the nanoflakes with four morphologies. The yellow arrows delineate the growth direction of γ -Fe₂O₃ nanoflakes. Figure 4e confirms that the triangular γ -Fe₂O₃ nanoflake grew along three diffusion pathways. The AFM images (Figures 4b and f) show the surface of the stellated nanoflake similar to two triangular Lego blocks rotationally aligned and stacked together. The yellow arrow marked in Figure 4f reveals a similar growth model compared with the triangular nanoflake. In accordance with the introduction in the previous chapter, the triangular γ -Fe₂O₃ nanoflake had an epitaxial relationship with the sapphire substrate. Hence, the stellated nanoflake can be understood as two triangular γ -Fe₂O₃ nanoflakes grown on a substrate similar to AA- and AB-stacking configurations. The petaloid γ -Fe₂O₃ nanoflake (Figure 4c) looked like a fresh blooming flower; its petals stretched out from the central cluster along the diffusion pathway (the yellow arrow marked in Figure 4g). Some upstretched “boundaries” existed among the petals. Figures S4a and b show the SEM images of two multipetaloid γ -Fe₂O₃ nanoflakes, which were grown following a typical diffusion-limited model. Similar petaloid graphene has also been reported, and this appearance is mostly caused by mutations in growth conditions.^{58, 63, 66} The chaotic growth state on the surface of dendritic γ -Fe₂O₃ nanoflakes (Figures 4d and f) indicated that the morphology of this nanoflake was difficult to control, although they exhibited sharp fan-shaped dendrites in the initial stage (Figures S4c and d). However, the dendritic γ -Fe₂O₃ nanoflake whose morphology is difficult to control is not the focus of this work.

On the basis of the descriptions of morphological characteristics and growth processes, we depicted a schematic (Figure 4i) for the four growth models. Such unique morphological characteristics were intuitively associated with the furnace temperature, that is, with the temperature of iron powder, the vapor pressure of iron in the tube was directly related to the temperature positively. The increase in temperature for the iron powder could cause the enlargement of nucleated clusters (Figures 4b and c) and the improvement of the growth rate of crystals. This viewpoint can be testified in the energy-dispersive spectroscopy maps shown in Figure S5. The iron content in the core of stellated and petaloid γ -Fe₂O₃ nanoflakes obtained at high temperature was considerably high. The size of the nanoflakes grown under high-temperature conditions was larger than others with the same growth time. The triangular γ -Fe₂O₃ flakes sequentially transform into stellated, petaloid, and dendritic crystalloids when enhancing the temperature of precursor powder and breaking the balance of the growth conditions (Figure 4i). This transition reveals the growth process from thermodynamically to kinetically dominated control and successfully realizes the morphology tunable synthesis recipe.^{56, 67}

The γ -Fe₂O₃ nanoflakes are air-stable materials that can be stored stably in the air for approximately 3 months (Figure S6).

Raman spectra and mapping images were collected to study the composition of the four shaped nanoflakes. Five uniform characteristic peaks exist in the Raman spectra of triangular, stellated, petaloid, and dendritic nanoflakes (Figure S7), proving the synthesis of γ -Fe₂O₃. Figure 5 provides the mapping images of Raman peak intensity, position, and width for triangular, stellated, and petaloid γ -Fe₂O₃ nanoflakes. The intensity, position, and width

mapping images (Figures 5a–c) of Raman A_{1g} peak for triangular $\gamma\text{-Fe}_2\text{O}_3$ nanoflakes exhibit a uniform signal distribution. As revealed in Figures 5d and g, the intensity signal of Raman A_{1g} peak for the centers of stellated and petaloid $\gamma\text{-Fe}_2\text{O}_3$ nanoflakes was considerably enhanced, which might be attributed to the existence of nucleated clusters. A uniform distribution of Raman A_{1g} peak in position and width signal was also observed (Figures 5e, f, h, and i). Figure 5g presents that the Raman signal at the “boundaries” of the petaloid $\gamma\text{-Fe}_2\text{O}_3$ nanoflake showed an unusual weak enhancement; thus, we collected some Raman spectra at these positions (spectra I–III in Figure S8). These Raman spectra could be deconvoluted to 14 Lorentz peaks; 3 of them (at ~ 419 , 579 , and 754 cm^{-1}) were contributed by sapphire. The extracted peak position, normalized intensity, and width of the 11 other peaks in Figure S8b are highly consistent with those of previously reported Raman curves for $\alpha\text{-Fe}_2\text{O}_3$.^{51, 68-70} A broad peak exists at $\sim 650\text{--}720\text{ cm}^{-1}$ in the Raman spectra of $\alpha\text{-Fe}_2\text{O}_3$ and $\gamma\text{-Fe}_2\text{O}_3$; consequently, the two phases with the mapping images of the peak at this position are difficult to distinguish. For this reason, the typical Raman peak ($\sim 1315\text{ cm}^{-1}$) of $\alpha\text{-Fe}_2\text{O}_3$ was selected to plot the intensity, position, and width mapping images (Figures 5j–l), which proved that $\alpha\text{-Fe}_2\text{O}_3$ appeared at the “boundaries” and core of the petaloid nanoflake. Raman spectra IV and V in Figure S8 were collected at the edge of dendritic nanoflakes. These observations indicated that the triangular, stellated, petaloid, and dendritic nanoflakes were all $\gamma\text{-Fe}_2\text{O}_3$, but the petaloid and dendritic nanoflakes contained a small amount of $\alpha\text{-Fe}_2\text{O}_3$.

To gain a deep insight into the oxidation state of iron, we conducted XPS measurements on the as-grown triangular, stellated, petaloid, and dendritic nanoflakes (Figure S9). Figures S9a and b are the high-resolution XPS spectra of Fe and O elements for triangular nanoflakes,

which have been introduced in the previous chapter. Figures S9c, e, and g show the chemical state of Fe element for stellated, petaloid, and dendritic nanoflakes, respectively. All the peaks were fitted using the Gauss–Lorentz (ratio: 40%) combination formula. The corresponding high-resolution XPS spectra of the O1s region (Figures S9d, f, and h) have only one peak, which was fitted using the Gauss–Lorentz (ratio: 30%) combination formula. Table S1 lists all the fitting results of the chemical state for Fe 2p and O 1s. The longitudinal comparison indicated that the position and full width at half maximum of each peak were homogeneous, indicating the purity of synthetic γ -Fe₂O₃ nanoflakes.

γ -Fe₂O₃ is a traditional intrinsic FM.^{35, 71-73} We performed magnetic hysteresis (*M-H*) loop measurements to verify the magnetic performance of the straticulate γ -Fe₂O₃ nanoflakes. Figures S10–13 show the normal state *M-H* loops of triangular, stellated, petaloid, and dendritic γ -Fe₂O₃ nanoflakes, indicating that these γ -Fe₂O₃ nanoflakes were ferromagnetic in the range of 5 K to 300 K. The diamagnetic contribution from the Al₂O₃ substrate was subtracted. To deepen the magnetic understanding of γ -Fe₂O₃ nanoflakes, the H_C , saturation magnetization (M_S), and remnant magnetization (M_R) in vertical (\perp , out-of-plane) and parallel (\parallel , in-plane) magnetic fields were also extracted and plotted. Figures 6a and b show the *M-H* loops at 5 K in vertical and parallel magnetic fields, respectively; Figures 6c and d are the *M-H* loops at 300 K in vertical and parallel magnetic fields, respectively. These figures demonstrate that the magnetic moment and saturation magnetization (M_S) of the triangular, stellated, petaloid, and dendritic γ -Fe₂O₃ nanoflakes increased in turn. The parallel magnetic field led to steep slopes for *M-H* loops in the near-zero magnetic field. The difficulty of magnetizing an FM in the direction of a crystal axis can be represented by the energy required to magnetize a unit FM to

saturation (E_M). The larger E_M is, the more difficult it is to magnetize, and vice versa. E_M can be expressed as follows:

$$E_M = \frac{W}{V} = \frac{1}{V} \int_0^{M_S} H dM,$$

where W is the work done in the magnetization process, which is also the area between the initial M - H curve and the M axis; V means the volume of the FM; M_S represents the saturation magnetization; and H represents the initial M - H curve. As shown in Figure S14, the work required for complete magnetization of petaloid γ - Fe_2O_3 nanoflakes in vertical and parallel magnetic fields at 5 K indicated that minimal work needs to be done in vertical and magnetic fields. Therefore, out-of-plane is the easy magnetization axis for γ - Fe_2O_3 nanoflakes, which is consistent with the fact that the M - H curves exhibit great H_C in vertical and magnetic fields. Figures 6e and f summarize the relationship between H_C and temperature of the triangular, stellated, petaloid, and dendritic γ - Fe_2O_3 nanoflakes in vertical and parallel magnetic fields, respectively. H_C increased with the decrease in temperature. Compared with the triangular and stellated γ - Fe_2O_3 nanoflakes, the petaloid and dendritic γ - Fe_2O_3 nanoflakes showed strong H_C ; this condition may be attributed to the existence of a small amount of α - Fe_2O_3 in the petaloid and dendritic γ - Fe_2O_3 nanoflakes, which corresponded to the growth process transition from thermodynamics to kinetics. In particular, α - Fe_2O_3 has exhibited ferromagnetism in some cases.⁷⁴⁻⁷⁶ This method of adjusting H_C in a large range by changing the shape of γ - Fe_2O_3 nanoflakes is an interesting event for the application of 2D γ - Fe_2O_3 nanoflakes.

We further investigated the magnetoelectric performance of the γ - Fe_2O_3 nanoflakes that were used in this study and synthesized at 930 °C and 980 °C; metal electrodes (Cr/Au) were then fabricated on top by using standard photoetching (Figures 7a and b). The electrical

transport data were from 400 K to low temperature were adopted. Figure 7c is the Hall resistance (R_H) response of a triangular γ -Fe₂O₃ device (Figures 7a and b) to a magnetic field sweep, which exhibited a feeble classical rectangular hysteresis loop at room temperature (300 K). As shown in Figure 7d, the electrical resistivity (ρ) of the triangular γ -Fe₂O₃ device rapidly increased monotonously with decreasing temperature in the whole temperature range, presenting typical semiconductive performance ($\rho = R_{xx}S/L$, where R_{xx} is the longitudinal resistance, S is the cross-sectional area, and L represents the length of the device).²⁶ These ρ - T curves are relatively close even if the device (Figure S15) thickness is different. Figure 7e exhibits that the ρ depended on the thickness of triangular and petaloid γ -Fe₂O₃ devices (Figure S16), indicating that the two morphologies of γ -Fe₂O₃ nanoflakes show consistent resistivity. In general, the triangular γ -Fe₂O₃ nanoflakes exhibit room-temperature intrinsic ferromagnetism and semiconductive performance, and their resistance can be effectively adjusted by changing the temperature. The γ -Fe₂O₃ nanoflakes with diverse shapes synthesized by controlling temperature have similar resistivity.

Conclusions

We achieved epitaxial growth of 2D nonlayered triangular γ -Fe₂O₃ nanoflakes on sapphire. Their lateral size is mainly affected by growth time. The γ -Fe₂O₃ nanoflakes can be controlled by increasing the precursor temperature to exhibit triangular, stellated, petaloid, and dendritic shapes in sequence, which represent a morphology tunable transition process from thermodynamically to kinetically dominated control. Although all the γ -Fe₂O₃ nanoflakes exhibit room-temperature ferromagnetism, the petaloid and dendritic γ -Fe₂O₃ nanoflakes show

strong H_C compared with the triangular and stellated $\gamma\text{-Fe}_2\text{O}_3$ nanoflakes. This difference may be attributed to the existence of a small amount of $\alpha\text{-Fe}_2\text{O}_3$ in the petaloid and dendritic $\gamma\text{-Fe}_2\text{O}_3$ nanoflakes. The triangular $\gamma\text{-Fe}_2\text{O}_3$ nanoflakes exhibit room-temperature intrinsic ferromagnetism and semiconductive performance, and their resistance can be effectively adjusted by changing the temperature. Various $\gamma\text{-Fe}_2\text{O}_3$ nanoflakes synthesized at different temperatures show room-temperature intrinsic ferromagnetism and have similar electrical resistivity. Our work provides a fundamental understanding of control synthesis and performance for 2D nonlayered $\gamma\text{-Fe}_2\text{O}_3$ nanoflakes. Overall, the synthesis of ultrathin, air-stable, and room-temperature ferromagnetic $\gamma\text{-Fe}_2\text{O}_3$ nanoflakes with different shapes and properties paves the way for the preparation and application of next generation electronic and spintronic devices.

METHODS

Growth of $\gamma\text{-Fe}_2\text{O}_3$ nanoflakes and Films. The $\gamma\text{-Fe}_2\text{O}_3$ nanoflakes were grown on $\alpha\text{-Al}_2\text{O}_3$ substrates in a home-made chemical vapor deposition (CVD) system (Figure S1a). A corundum boat containing mixed powder of 0.035 g Fe (99.9%, Alfa Aesar) and 0.01 g Co (99.8%, Alfa Aesar) was placed in the temperature zone, and the cleaned $\alpha\text{-Al}_2\text{O}_3$ (*c*-plane (0001)) substrate was face-down covered on the corundum boat. Before heating, the CVD system was evacuated and then flushed with 200 sccm of high-purity argon gas for 30 min to remove air and other impurities. After purging the system, the furnace was heated to 930~1080°C. In the meanwhile, 50 sccm argon gas was introduced in the CVD system as carrier gas under atmospheric pressure. The reaction time ranged from 3 to 30 min, and the high-temperature system was cooled down

naturally to room temperature. The γ -Fe₂O₃ nanoflakes or films were eventually obtained.

Structural characterization. OM images were taken using an optical microscope (PSM-1000, Moti) equipped with three long focal length lenses ($\times 10$, $\times 50$, and $\times 100$). Raman spectra were obtained using a confocal Raman microscope (Alpha 300R, WITec) with a 532 nm laser, and the laser power was measured by using a hand-held optical power meter console (PM100D, Thorlabs). The laser spot diameter was 1 μm . XRD analyses were performed using an X-Ray Diffractometer (X PERT PRO MRD S4, PANalytical). XPS measurements were conducted on a spectrometer (ESCALAB 250Xi, Thermo Scientific) with Al K_{α} X-rays (1489.6 eV). SEM images were taken using a field-emission scanning electron microscope (QUANTA 650FEG, FEI). AFM images were obtained using an atomic force microscope (Icon, Bruker). TEM and HRTEM observations were carried out by a transmission electron microscope (JEM-2100, JEOL). The TEM specimen was prepared by the standard wet transfer method.

Device Fabrication and property measurement. The Hall patterns were fabricated using the standard photolithography with the help of direct-write laser-based micro-fabrication system (DWL 2000, Heidelberg Instruments) and Photoresist Coater (Gamma Cluster System, Karls SUSS MicroTec). The Cr/Au (10/70 nm) metal electrodes were then sputtered on top using a multitarget confocal magnetron sputter (UHV PVD, Kenosistec). Magnetization measurements were performed by using a commercial superconducting quantum interference device (SQUID) (MOMS 3, QUANTUM DESIGN) with the magnetic field either vertical (out-of-plane) or parallel (in-plane) to the nanoflakes. The Al₂O₃ substrates covered by γ -Fe₂O₃ nanoflakes were diced to $3 \times 5 \text{ mm}^2$ chips and loaded into a hollow copper rod. The magnetic hysteresis loops were measured from 5 K to 300 K while cycling a DC magnetic field through -2 T to 2 T. The

magnetoelectric data were collected by using a Hall measurement system (HMS 9709A, Lake Shore), where magnetic field was applied parallel to the nanoflakes in measurement.

Author Contributions

Z.C.W., W.J.W., and Y.L. initiated and managed the project. Z.Y.J., W.J.W., and Z.C.W. completed the sample growth, characterization, device fabrication and prepared the manuscript. Z.C.L. and W.J.W. performed the magnetic and magnetoelectric measurements. All authors discussed the results and directed the entire study.

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Figure 1

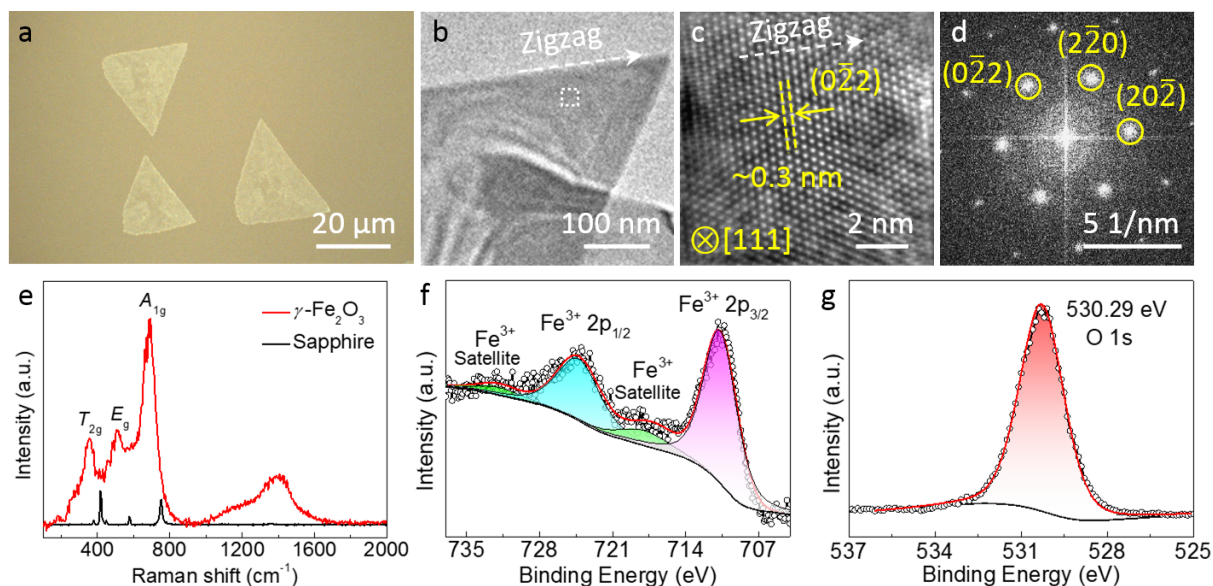


Figure 1 Characterization of γ - Fe_2O_3 nanoflakes. (a) Optical image of three triangular γ - Fe_2O_3 nanoflakes grown on sapphire substrate. (b) TEM image of a transferred γ - Fe_2O_3 nanoflake. (c) High-resolution TEM image taken from the square-marked area in (b), which is projected along $[111]$ zone axis of the γ - Fe_2O_3 nanoflake. The white arrow indicates the zigzag direction, as also marked in a white arrow in (b). The $(0\bar{2}2)$ plane spacing is ~ 0.3 nm. (d) Corresponding FFT patterns of the HRTEM image in (c), which displays only one set of sixfold symmetry spots. (e) Raman spectra of the γ - Fe_2O_3 nanoflake and sapphire. The red curve contains five main peaks at ~ 356 , 512 , 650 – 720 , 1200 , and 1400 cm^{-1} . (f,g) High-resolution XPS spectra of the Fe 2p (f) and O 1s (g) region.

Figure 2

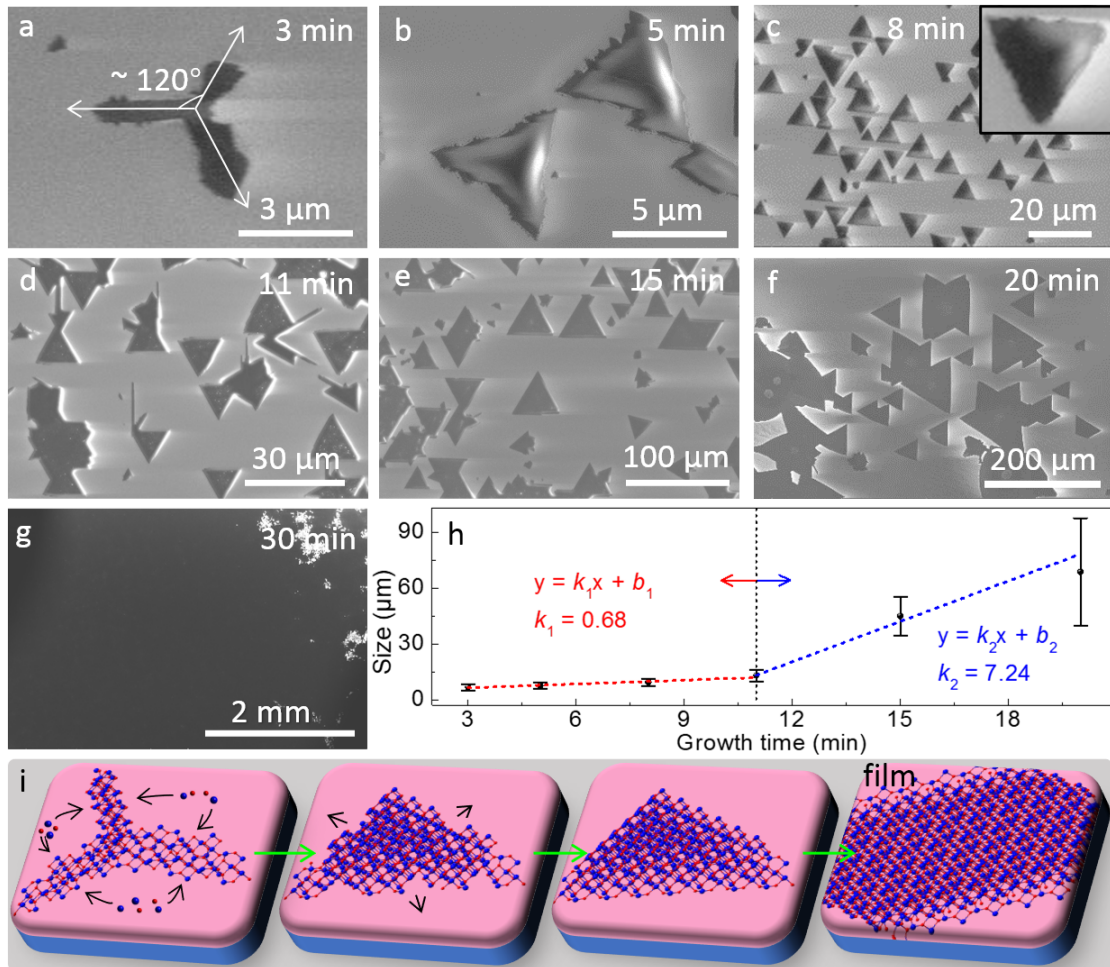


Figure 2 Characterization of the growth process of triangular γ -Fe₂O₃ nanoflakes. (a–g) SEM images of the as-grown γ -Fe₂O₃ nanoflakes with the increase of growth time (3, 5, 8, 11, 15, 20, and 30 min), revealing the effect of growth time on crystal shape and size. (h) Size of γ -Fe₂O₃ nanoflakes as a function of growth time, showing linear relationships between size and growth time at different stages. The slopes of the initial and rapid growth stages are estimated to be 0.68 and 7.24, respectively. (i) Schematic of the growth process of γ -Fe₂O₃, ranging from dendrite formation, filling defects, growth, to formation of film.

Figure 3

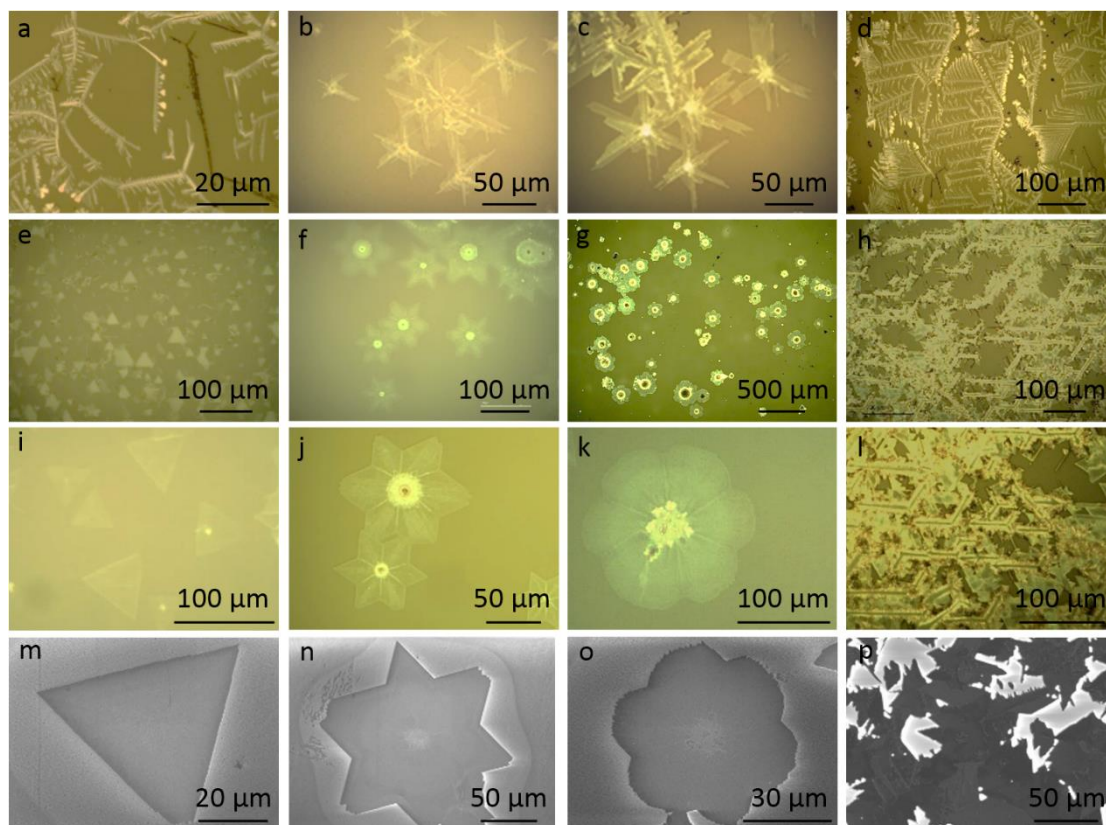


Figure 3 Characterization of the stage-by-stage formation of the γ -Fe₂O₃ nanoflakes grown at different temperatures (\sim 930°C, 980°C, 1030°C, and 1080°C), uncovering five growth models. (a–d) OM images of the as-grown triangular (a), stellated (b), petaloid (c), and dendritic (d) γ -Fe₂O₃ nanoflakes at the initial growth stage. (e–h) OM images of the as-grown triangular (e), stellated (f), petaloid (g), and dendritic (h) γ -Fe₂O₃ nanoflakes at the mature growth stage. (i–l) High-magnification OM images corresponding to the low-magnification OM images in (e–h). (m–p) SEM images of triangular (m), stellated (n), petaloid (o), and dendritic (p) γ -Fe₂O₃ nanoflakes.

Figure 4

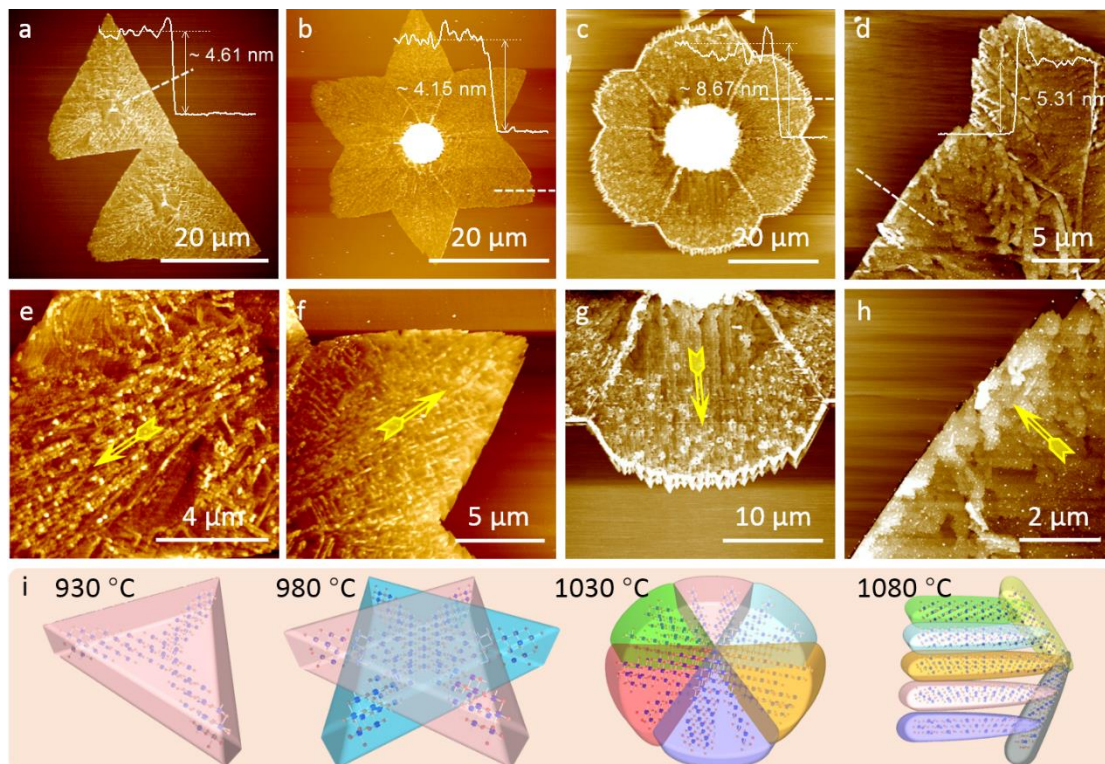


Figure 4 Characterization of surface profile. (a–d) AFM images of triangular (a), stellated (b), petaloid (c), and dendritic (d) γ -Fe₂O₃ nanoflakes. The insets show the corresponding profile lines of the γ -Fe₂O₃ nanoflakes, indicating that the thickness of the nanoflakes in (a) to (d) is about 4.61, 4.15, 8.67, and 5.31 nm, respectively. (e–h) High-magnification AFM images corresponding to the low-magnification AFM images in (a–d). The yellow arrows delineate the growth pathways of γ -Fe₂O₃ nanoflakes. (i) Sketch of the triangular, stellated, petaloid, and dendritic γ -Fe₂O₃ nanoflakes synthesized at 930°C, 980°C, 1030°C, and 1080°C.

Figure 5

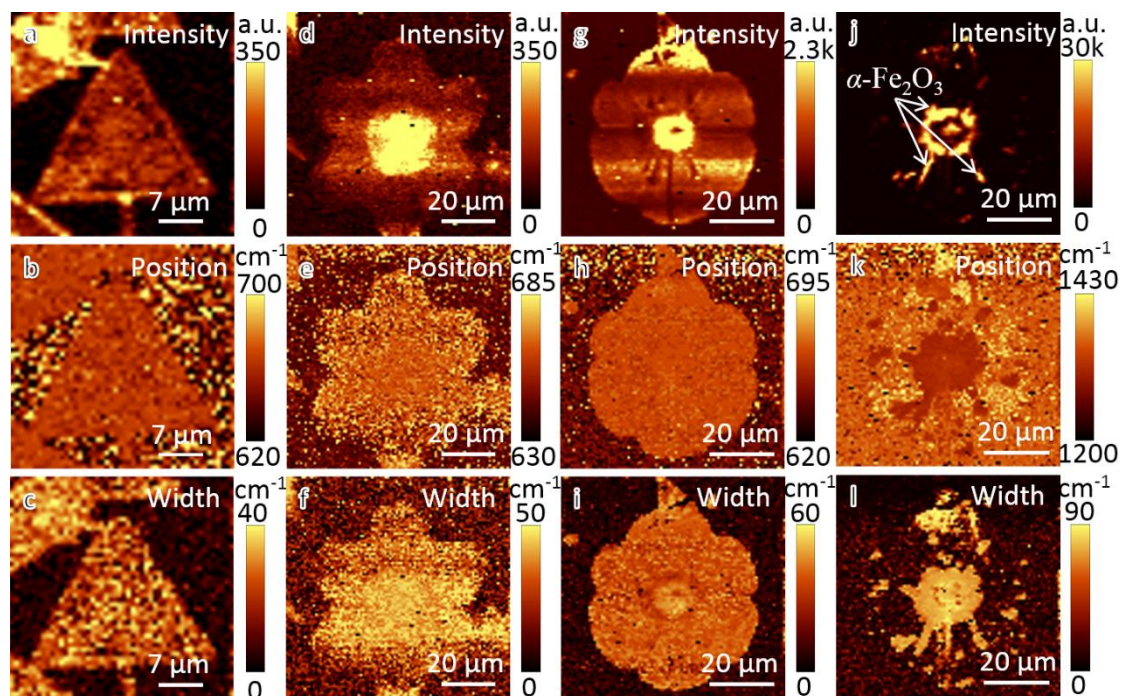


Figure 5 Raman characterization of the triangular, stellated, petaloid, and dendritic γ - Fe_2O_3 nanoflakes. (a–c) Raman A_{1g} mode mapping images of intensity (a), position (b), and width (c) for a triangular γ - Fe_2O_3 nanoflake. (d–f) Raman A_{1g} mode mapping images of intensity (d), position (e), and width (f) for a stellated γ - Fe_2O_3 nanoflake. (g–i) Raman A_{1g} mode mapping images of intensity (g), position (h), and width (i) for a petaloid γ - Fe_2O_3 nanoflake. (j–l) Intensity (j), position (k), and width (l) mapping images of a Raman peak at $\sim 1315\text{ cm}^{-1}$ for the petaloid γ - Fe_2O_3 nanoflake in (g–i).

Figure 6

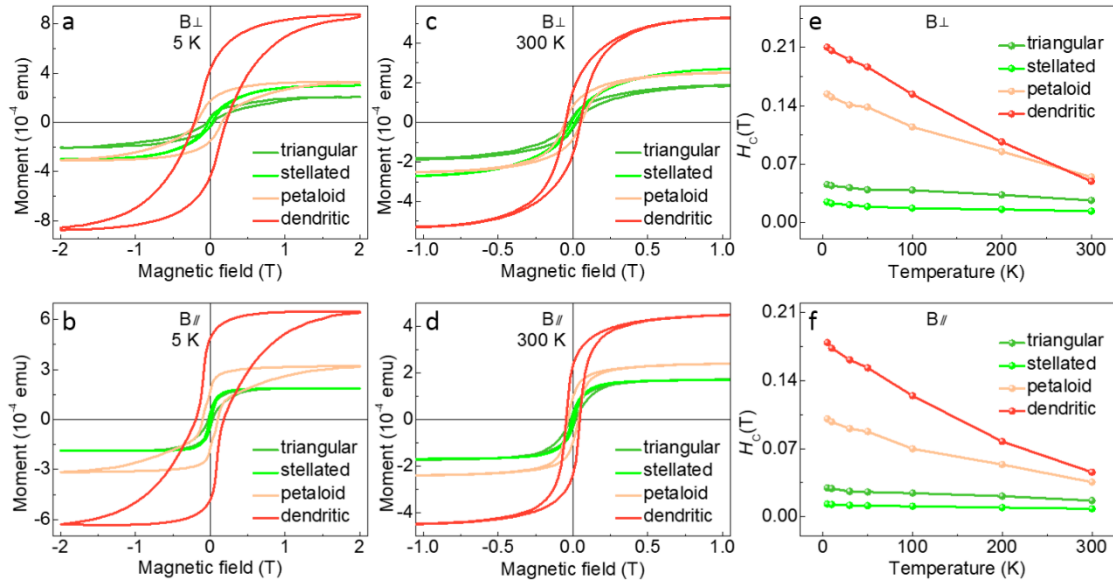


Figure 6 Characterization of ferromagnetic performance of $\gamma\text{-Fe}_2\text{O}_3$ nanoflakes. (a,b) Magnetic hysteresis ($M-H$) loops of the triangular, stellated, petaloid, and dendritic $\gamma\text{-Fe}_2\text{O}_3$ nanoflakes measured under vertical (a) and parallel (b) fields at 5 K. (c,d) $M-H$ loops of triangular, stellated, petaloid, and dendritic $\gamma\text{-Fe}_2\text{O}_3$ nanoflakes measured under vertical (c) and parallel (d) fields at 300 K. The diamagnetic contribution from $\alpha\text{-Al}_2\text{O}_3$ substrate is subtracted. (e,f) Temperature-dependent coercive force (H_c) for $\gamma\text{-Fe}_2\text{O}_3$ nanoflakes measured under vertical (e) and parallel (f) magnetic fields.

Figure 7

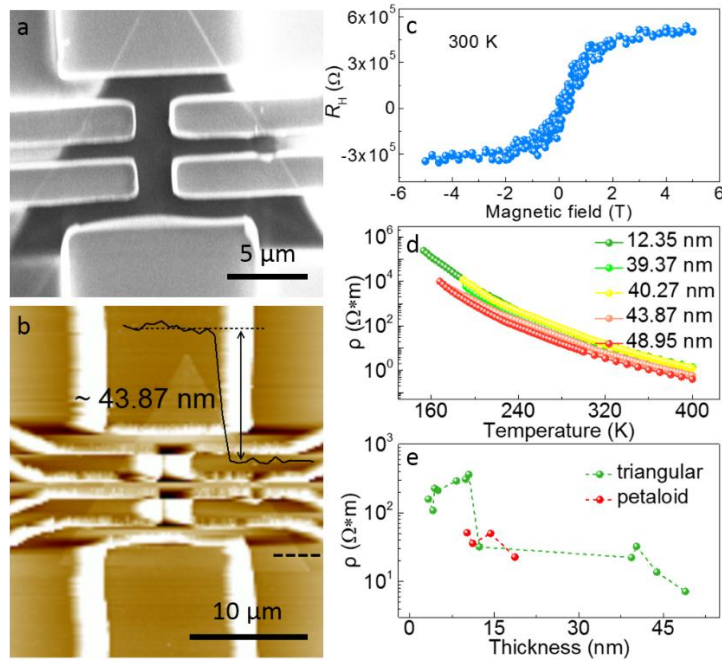


Figure 7 Magnetic and electrical transport measurements of $\gamma\text{-Fe}_2\text{O}_3$ nanoflakes. (a,b) SEM (a) and AFM (b) images of a triangular $\gamma\text{-Fe}_2\text{O}_3$ Hall device. The inset shows the height profile line of the Hall device. (c) Hall resistance (R_H) as a function of magnetic field measured at 300 K for the Hall device shown in (a). (d) Temperature-dependent electrical transport measurement of the electrical resistivity (ρ) as thickness varies (12.35, 39.37, 40.27, 43.87, and 48.95 nm). (e) Thickness-dependent ρ for triangular and petaloid $\gamma\text{-Fe}_2\text{O}_3$ nanoflakes.